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EXPONENTIAL STABILITY OF POSITIVE SOLUTIONS TO SOME NONLINEAR HEAT EQUATIONS

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Abstract: Following a recent work of A. Haraux in which he proves exponential stability of positive solutions of a heat equation with strictly convex nonlinearity, the same property is shown for a suitable perturbation of the nonlinearity which can, in particular, be non convex.

1 – Introduction and main results

Let Ω be a bounded and connected open subset of \mathbb{R}^N with a Lipschitz continuous boundary and let us consider the semilinear heat equation

(1.1)
$$u_t - \Delta u + f(u) = k(t, x) \quad \text{in } \mathbb{R}^+ \times \Omega$$
$$u(t, \cdot) = 0 \quad \text{on } \mathbb{R}^+ \times \partial \Omega ,$$
$$u(0, \cdot) = u_0(\cdot) \quad \text{in } \Omega ,$$

and the elliptic equation

(1.2)
$$\begin{aligned} -\Delta u + f(u) &= 0 \quad \text{in } \Omega ,\\ u &= 0 \quad \text{on } \partial \Omega , \end{aligned}$$

where $f: \mathbb{R} \to \mathbb{R}$ is a locally Lipschitz continuous function such that

(1.3)
$$f(0) = 0$$
 and $f(s) \to +\infty$ as $s \to +\infty$

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and $k \colon \mathbb{R}^+ \times \Omega \to \mathbb{R}$ satisfies the conditions

(1.4)
$$k \in L^{\infty}(\mathbb{R}^+ \times \Omega)$$
 and $k(t, x) \ge 0$ a.e. on $\mathbb{R}^+ \times \Omega$.

By using standard techniques from the theory of evolution equations, cf. e.g. [5], we know that for all $u_0 \in L^{\infty}(\Omega)$ with $u_0(x) \ge 0$ a.e. on Ω , there exists a unique solution $u \in C((0, +\infty); H_0^1(\Omega) \cap L^{\infty}(\Omega)) \cap C([0, +\infty); L^2(\Omega))$ of (1.1) such that $u(0, \cdot) = u_0(\cdot)$. In addition we have

$$u(t,x) \ge 0$$
 a.e. on $\mathbb{R}^+ \times \Omega$.

As a consequence of (1.3) and the maximum principle, u is uniformly bounded on $\Omega \times \mathbb{R}^+$. Then by the method of [11], it follows easily that

$$\bigcup_{t\geq 1} \{u(t,\cdot)\} \quad \text{ is bounded in } \ C^{1+\alpha}(\overline{\Omega}) \ \text{ for every } \alpha \in [0,1) \ .$$

In particular the curve $t \mapsto u(t, \cdot)$ has a precompact range in $H_0^1(\Omega) \cap L^{\infty}(\Omega)$ for $t \geq 1$ and it is natural to ask about the asymptotic behavior of $u(t, \cdot)$ as $t \to \infty$.

A. Haraux [8] has proved exponential convergence of nonnegative solutions of (1.1) when f satisfies the additional hypotheses

(1.5)
$$f$$
 strictly convex on $[0, +\infty)$ and $f'_d(0) < -\lambda_1(-\Delta)$

where $\lambda_1(-\Delta)$ is the smallest eigenvalue of $(-\Delta)$ in $H_0^1(\Omega)$. The proof of this result is based on the uniqueness of positive solution of the equation (1.2) and the fact that $\lambda_1(-\Delta + f'(\varphi)) > 0$ (φ is the unique positive solution of (1.2)).

The typical example of nonlinearities which verifies these hypotheses is the following

(1.6)
$$f(s) = s^p - \lambda s, \quad \lambda > \lambda_1(-\Delta), \quad p > 1.$$

The question which we study in this paper is the following: What happens if we perturb the nonlinearity in such a way that convexity of f is lost? In the special case of example (1.6) a question of interest is the following: Can we find $\varepsilon > 0$ such that the result of [8] persists for the new nonlinearity

$$h(s) = s^p - \lambda s - \varepsilon s^q$$

with p, λ as in (1.6) and 1 < q < p?

We are able to give a positive answer to this question. We use the same method as in [8]: At first we prove the uniqueness of positive solution of (1.2)

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with this new type of nonlinearity. We assume the following hypotheses: Let f satisfying (1.3), (1.5), and let $g: \mathbb{R}^+ \to \mathbb{R}$ be a function of class C^1 such that

(1.7)
$$g'(0) \ge 0, \quad \lim_{s \to \infty} f(s) - C g(s) = \infty, \\ g(0) = 0, \quad g(s) \ge 0 \quad \forall s \ge 0,$$

with C > 0 and we consider the nonlinear heat equation

(1.8)
$$u_t - \Delta u + f(u) = \varepsilon g(u) + k(t, x) \quad \text{in } \mathbb{R}^+ \times \Omega ,$$
$$u(t, \cdot) = 0 \quad \text{on } \mathbb{R}^+ \times \partial \Omega ,$$
$$u(0, \cdot) = u_0(\cdot) \quad \text{in } \Omega$$

The main results of this paper are the following

Theorem 1.1. Let f, g satisfy the hypotheses (1.3), (1.5), (1.7). Then there exists $\varepsilon_1 > 0$ such that for all $\varepsilon \in [0, \varepsilon_1)$ the equation

(1.9)
$$\Psi \in H_0^1(\Omega) , \quad -\Delta \Psi + f(\Psi) = \varepsilon g(\Psi) ,$$

has one and only one solution $\Psi \ge 0$ other than 0. In addition we have $\Psi > 0$ everywhere in Ω and

(1.10)
$$\lambda_1 \Big(-\Delta + f'(\Psi) - \varepsilon g'(\Psi) \Big) > 0 \quad \forall \varepsilon \in [0, \varepsilon_1) .$$

Theorem 1.2. Let f, g and k satisfy the hypotheses (1.3), (1.4), (1.5), (1.7). Then if $u_0, v_0 \in L^{\infty}$ with $u_0(x) \ge 0$ and $v_0(x) \ge 0$ a.e. on Ω , consider the solution u, v of (1.1) with respective initial data $u(0, x) = u_0(x)$ and $v(0, x) = v_0(x)$. Assuming either that both u_0, v_0 are not identically 0 or that k(t, x) > 0 on a subset of positive measure of $\mathbb{R}^+ \times \Omega$. Then there exists $\varepsilon_2 > 0$ such that for all $\varepsilon \in [0, \varepsilon_2)$, there is $\gamma > 0$ independent of k and (u_0, v_0) :

(1.11)
$$\forall t \ge 0 \quad \|u(t, \cdot) - v(t, \cdot)\|_{\infty} \le C(u_0, v_0, \varepsilon) \exp(-\gamma t) .$$

The paper is organized as follows: in Section 2 we prove Theorem 1.1, in Section 3 we establish Theorem 1.2 when k = 0. In Section 4, we establish Theorem 1.2 in the general case. In each section some remarks are presented.

2 – The stationnary problem

The object of this section is to prove Theorem 1.1.

Proof of Theorem 1.1. First we prove the existence of a positive solution for the equation (1.9). In fact, if $\varepsilon = 0$ then by a theorem of Berestycki [1] (Theorem 4, page 14, cf. also [2], [3]), there exists a unique positive solution φ of (1.2) which verifies

(2.1)
$$\lambda_1(-\Delta + f'(\varphi) \, id) > 0$$

Since $g \ge 0$ then φ is a subsolution of (1.9).

Now we assume that $\varepsilon < C$, then by (1.7) there exists M > 0 such that

(2.2)
$$f(M) - \varepsilon g(M) > 0 .$$

So M is a supersolution of (1.9). We claim that $\|\varphi\|_{\infty} < M$. Indeed, let $x_0 \in \Omega$ such that $\varphi(x_0) = \|\varphi\|_{\infty}$, we have $\Delta\varphi(x_0) \leq 0$. Now if $\|\varphi\|_{\infty} \geq M$ then we have $f(M) \leq 0$. Hence $f(M) - \varepsilon g(M) \leq 0$, and this contradicts (2.2). Then there exist a solution Ψ for (1.9) which verifies $\varphi \leq \Psi \leq M$. By using again the maximum principle, we have for all ξ positive solution of (1.9) $\xi < M$. Then the problem (1.9) has a "maximal" solution Ψ in the sense: any solution $\xi \neq \Psi$ of (1.9) is less than Ψ . (This solution can be constructed by a standard iterative scheme.)

Now we have to use the following lemma due to Haraux [9, 10].

Lemma 2.1. Let f satisfy the hypotheses (1.3), (1.5) and let φ be the positive solution of the equation

$$\varphi \in C(\Omega) \cap H_0^1(\Omega), \quad -\Delta \varphi + f(\varphi) = 0.$$

Let on the other hand $\xi \geq 0$ be a solution of

$$\xi \in C(\Omega) \cap H_0^1(\Omega), \quad -\Delta \xi + f(\xi) \ge_0.$$

Then either $\xi = 0$ or $\xi \ge \varphi$.

Proof of Theorem 1.1 (continued). Let ξ be a positive solution of (1.9), then by using Lemma 2.1 we have

$$\varphi \leq \xi < M \ .$$

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Now we prove uniqueness. In fact we assume that we have a solution ξ of (1.9) other than the "maximal" solution Ψ . Then we have:

(2.3)
$$-\Delta(\Psi-\xi) + f(\Psi) - f(\xi) = \varepsilon[g(\Psi) - g(\xi)] .$$

Multiplying (2.3) by $(\Psi - \xi)$ and integrating over Ω we find

(2.4)
$$\int_{\Omega} |\nabla(\Psi - \xi)|^2 + [f(\Psi) - f(\xi)] (\Psi - \xi) \, dx = \\ = \varepsilon \int_{\Omega} [g(\Psi) - g(\xi)] (\Psi - \xi) \, dx \, .$$

Since $\varphi \leq \xi \leq \Psi < M$, then by using (1.5), (1.7) and (2.4) we find

(2.5)
$$\int_{\Omega} |\nabla(\Psi - \xi)|^2 + f'(\varphi) |\Psi - \xi|^2 dx \le \varepsilon \int_{\Omega} C_1 |\Psi - \xi|^2 dx$$

with $C_1 = \sup\{|g'(s)|, s \in [0, M]\} > 0$. So

(2.6)
$$\left[\lambda_1(-\Delta + f'(\varphi)) - \varepsilon C_1\right] \int_{\Omega} |\Psi - \xi|^2 \, dx \le 0 \; .$$

Thank's to (2.1) $\lambda_1(-\Delta + f'(\varphi) id) > 0$. Now let ε' such that $\lambda_1(-\Delta + f'(\varphi) id) = \varepsilon' C_1$ and $\varepsilon_1 = \inf(\varepsilon', C)$, with C as in (1.7). Then for all $\varepsilon \in [0, \varepsilon_1)$, we have

(2.7)
$$\lambda_1(-\Delta + f'(\varphi)) - \varepsilon C_1 > 0 .$$

The uniqueness follows from (2.7), we note this solution by Ψ . By using (1.5), (1.7) and (2.7) we deduce

(2.8)
$$\lambda_1 \Big(-\Delta + [f'(\Psi) - \varepsilon g'(\Psi)] \, id \Big) > 0 \quad \forall \varepsilon \in [0, \varepsilon_1) \, . \blacksquare$$

3 - The autonomous case

The object of this section is to prove Theorem 1.2 in the case k = 0. We use the method of [8].

Proof of Theorem 1.2. Let $Z = \{z \in C(\overline{\Omega}) \cap H_0^1(\Omega) \mid z \ge 0\}$. Subsequently $h = f - \varepsilon g$ with $\varepsilon \in [0, \varepsilon_1)$ and ε_1 is as in Theorem 1.1.

The equation (1.1) generates a dynamical system $\{S(t)\}_{t\geq 0}$ which assigns to each element $z \in Z$ the value v(t) = S(t) z where v is the solution of (1.8) such that v(0) = z. Now let E be the functional defined by

$$\forall \varphi \in Z \quad E(\varphi) = \frac{1}{2} \int_{\Omega} |\nabla \varphi|^2 \, dx + \int_{\Omega} H(\varphi) \, dx \quad \text{where } H(u) \mathrel{\mathop:}= \int_0^u h(s) \, ds \; .$$

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E is a strict Liapunov functional on Z relative to S(t) and we refer to [9] for a simple proof.

Let $u_0 \in L^{\infty}(\Omega)$, $u_0 \geq 0$, then by using the maximum principle (cf. for example [5]) we have $u(t, x) \geq 0$ a.e. $(t, x) \in \mathbb{R}^+ \times \Omega$. By the standard invariance principle (cf. [9]), we conclude that the solution $u(t, \cdot)$ asymptotes the set of nonnegative solutions of (1.9) as $t \to \infty$. We now show that if $u_0 \neq 0$, $u(t, \cdot)$ cannot tend to 0 as $t \to \infty$.

In fact assuming that $\lim_{t\to\infty} ||u(t,\cdot)||_{\infty} = 0$, then for each $\alpha > 0$, there is $T(\alpha)$ such that

$$\forall t \ge T(\alpha) \quad h(u(t,x)) \le \{h'_d(0) + \alpha\} u(t,x) \quad \text{on } \Omega .$$

Choosing $\alpha > 0$ small enough such that $-h'_d(0) - \alpha - \lambda_1(\Omega) > 0$, multiplying the equation by the positive eigenfunction φ_1 corresponding to the first eigenvalue $\lambda_1(-\Delta)$ of $-\Delta$ in $H_0^1(\Omega)$ and integrating over Ω we find

$$\frac{d}{dt} \int_{\Omega} u(t,x) \varphi_1 \, dx \ge 0 \quad \forall t \ge T(\alpha) \; .$$

Since the function $t \mapsto \int_{\Omega} u(t, x) \varphi_1 dx$ is nondecreasing on $[T(\alpha), \infty]$ and tends to 0 as $t \to \infty$, it must vanish identically on $[T(\alpha), \infty]$. Because φ_1 is positive in Ω , this imply that $u(t, \cdot) = 0 \ \forall t \ge T(\alpha)$. Then a classical connectedness argument shows that $u_0 = 0$. Therefore if $u_0 \ne 0$, the ω -limit set of u_0 under S(t) is reduced to a single point: $\omega(u_0) = \{\Psi\}$. Since $u(t, \cdot)$ remain bounded in $C^1(\Omega)$ for all $t \ge 1$ we deduce that

$$\lim_{t\to\infty} \|u(t,\cdot) - \Psi(\cdot)\|_{1,\infty} = 0.$$

For the end of the proof, we just need to use (2.8).

Remark 3.1. It is clear that $\lim_{t\to\infty} ||u(t,\cdot) - \Psi(\cdot)||_{1,\infty} \exp(ct) = 0$, $\forall c < \lambda_1(-\Delta + h'(\psi) id)$. In [19], Wiegner has proved that in such a case the difference of two solutions tend to 0 as $\exp(-c_1 t)$ with $c_1 = \lambda_1(-\Delta + h'(\psi) id)$. For related works in the asymptotic of autonomous parabolic equation we refer to [7–9, 12–19].

4 – The nonautonomous case

The object of this section is to prove Theorem 1.2 in the general case. Subsequently $\varepsilon \in [0, \varepsilon_1)$ and ε_1 as in Theorem 1.1. In the proof we can use the following lemmas from [8] which are also valid for the modified equation (1.8):

Lemma 4.1. Let ψ be the unique positive solution of (1.9) and let us consider the solution z of (1.8) with initial condition $z(0) = \psi$. Then we have:

$$\forall t \ge 0 \quad z(t,x) \ge \psi(x) \text{ on } \Omega$$

Lemma 4.2. Let $u_0 \in L^{\infty}(\Omega)$ with $u_0(x) \geq 0$ a.e. on Ω and consider the solution u of (1.8) with initial datum $u(0, x) = u_0(x)$. Assuming either that u_0 is not identically 0 or that k(t, x) > 0 on a subset of positive measure of $\mathbb{R}^+ \times \Omega$, we have

(3.1)
$$\lim_{t \to \infty} \left\| \left(u(t, \cdot) - \psi(\cdot) \right)^{-} \right\|_{\infty} = 0 \; .$$

Proof of Theorem 1.2: Obviously, it is sufficient to prove the result when $v_0 = \psi$. Then v(t) = z(t) and

$$\begin{aligned} \forall t > 0 \quad & \frac{1}{2} \frac{d}{dt} \Big(\int_{\Omega} |u(t,x) - z(t,x)|^2 \, dx \Big) = \\ & = -\int_{\Omega} |\nabla(u-z)|^2 \, dx - \int_{\Omega} [f(u) - f(z)] \, (u-z) \, dx + \varepsilon \int_{\Omega} [g(u) - g(z)] \, (u-z) \, dx \; . \end{aligned}$$

By convexity of f, since $z(t) \ge \psi$ for all t, we have $f(z)/z \ge f(\psi)/\psi$. Moreover from (3.1) it follows in particular that fixing some nonempty open set ω contained in a compact subset of Ω , we have for $t \ge T$ depending on the solution u that

(3.2)
$$\forall t \ge T \quad u(t,x) \ge \frac{1}{2} \psi(x) \text{ on } \omega$$

Now from (3.2) we deduce easily the inequality

$$\begin{aligned} \forall t \ge T \quad &\frac{1}{2} \frac{d}{dt} \Big(\int_{\Omega} |u(t,x) - z(t,x)|^2 \, dx \Big) = \\ &= -\int_{\Omega} |\nabla(u-z)|^2 \, dx - \int_{\Omega} c(x) \, |u-z|^2 \, dx + \varepsilon \int_{\Omega} [g(u) - g(z)] \, (u-z) \, dx \end{aligned}$$

with

$$c(x) = \begin{cases} \frac{f(\psi)}{\psi} & \text{outside } \omega, \\ 2\frac{f(\psi) - f(\psi/2)}{\psi} & \text{in } \omega . \end{cases}$$

Let

$$\delta = \inf\left\{\int_{\Omega} \left(|\nabla w|^2 + c(x) w^2\right) dx, \ w \in H^1_0(\Omega), \ \int_{\Omega} w^2 dx = 1\right\}.$$

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We can prove as in [8] that $\delta > 0$. In the other hand, there exists $C_1 > 0$ such that

$$\int_{\Omega} [g(u) - g(z)] (u - z) \, dx \le C_1 \int_{\Omega} |u - z|^2 \, dx \; .$$

Set $\varepsilon'' = \frac{\delta}{C_1}$ and let $\varepsilon_2 = \inf(\varepsilon_1, \varepsilon'')$, then we obtain for all $t \ge T$

$$\frac{d}{dt} \left(\int_{\Omega} |u(t,x) - z(t,x)|^2 \, dx \right) \le -(\delta - \varepsilon \, C_1) \, \int_{\Omega} |u - z|^2 \, dx \, dx$$

The end of the proof is the same as in [8].

Remark 4.3. It is instructive to compare the result of Theorem 1.2 with the result of Chen and Matano [6], recently completed with a simple proof by Brunovsky et al. [4]. The result of [4, 6] are proved for any nonlinearity but only in one space dimension and for time-periodic forcing terms. On the other hand Theorem 1.2 is valid for any space dimension, but it is restricted to positive solution and a special type of nonlinearities.

Remark 4.4. This result can be viewed as a "structural stability" property for the result of [8]. However our method of proof is constructive since given $\lambda_1(-\Delta + f'(\psi) id) = \gamma > 0$, we can specify explicitly ε_1 and ε_2 in terms of the function g.

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