THE COMPRESSION OF A SLANT HANKEL OPERATOR TO H^2

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ABSTRACT. A slant Hankel operator K_{φ} with symbol φ in $L^{\infty}(T)$ (in short L^{∞}), where T is the unit circle on the complex plane, is an operator whose representing matrix $M=(a_{ij})$ is given by $a_{i,j}=\langle \varphi,z^{-2i-j}\rangle$, where $\langle \cdot,\cdot \rangle$ is the usual inner product in $L^2(T)$ (in short L^2). The operator L_{φ} denotes the compression of K_{φ} to $H^2(T)$ (in short H^2). We prove that an operator L on H^2 is the compression of a slant Hankel operator to H^2 if and only if $U*L=LU^2$, where U is the unilateral shift. Moreover, we show that a hyponormal L_{φ} is necessarily normal and L_{φ} can not be an isometry.

1. Introduction

Let φ be in L^{∞} . Then $\varphi(z) \sim \sum_{i=-\infty}^{\infty} a_i z^i$, where $a_i = \langle \varphi, z^i \rangle$ is the *i*-th Fourier coefficient of φ and $\{z^i : i \in Z\}$ is the usual orthonormal basis of L^2 and Z is the set of integers. A slant Toeplitz operator A_{φ} is an operator on L^2 defined by

$$A_{\varphi}(z^k) = \sum_{i=-\infty}^{\infty} a_{2i-k} z^i,$$

for k in Z. Furthermore $A_{\varphi} = WM_{\varphi}$, where M_{φ} is a multiplication operator on L^2 and W is an operator on L^2 such that $Wz^{2n} = z^n$ and $Wz^{2n-1} = 0$, for n in Z.

A Hankel operator S_{φ} is an operator on L^2 defined by

$$S_{\varphi}(z^k) = \sum_{i=-\infty}^{\infty} a_{-i-k} z^i$$

for k in Z [1]. Moreover, $S_{\varphi} = JM_{\varphi}$ and $M_{\varphi} = JS_{\varphi}$, where J is the reflect in operator on L^2 , that is, $J(z^n) = z^{-n}$, for n in Z. A slant Hankel operator K_{φ} is

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an operator on L^2 defined by

$$K_{\varphi}(z^k) = \sum_{i=-\infty}^{\infty} a_{-2i-k} z^i$$

for k in Z. Moreover, $K_{\varphi} = JA_{\varphi}$ and $A_{\varphi} = JK_{\varphi}$, The compression of K_{φ} to H^2 is denoted by L_{φ} . Symbolically $L_{\varphi} = PK_{\varphi}|_{H^2}$, equivalently $L_{\varphi}P = PK_{\varphi}P$, where P is the orthogonal projection of L^2 onto H^2 .

In this paper, we establish equations which characterize slant Hankel operators and their compressions. We also investigate the conditions under which these operators are self-adjoint, normal, hyponormal or compact.

2. Slant Hankel operators

In this section, we obtain a characterization of a slant Hankel operator on L^2 and prove that the set of all slant Hankel operators on L^2 is a subspace of $\mathbf{B}(L^2)$ (the algebra of all operators on L^2). We begin by recalling the definition of the reflection operator J on L^2 . For f in L^2 , $J(f(z)) = f(\bar{z})$. The matrix of J with respect to the orthonormal basis $\{z^n : n \in Z\}$ is

Since $J^2(f(z)) = J(f(\bar{z})) = f(z)$, it follows that $J^2 = I$. Since $\langle J^*(f(z)), g(z) \rangle = \langle f(z), J(g(z)) \rangle = \langle f(z), g(\bar{z}) \rangle = \langle f(\bar{z}), g(z) \rangle$, for all f, g in L^2 , it follows that $J^* = J$. Moreover, we have ||J|| = 1.

REMARK 2.1. (a) Since, for n in Z, $WJ(z^{2n})=\bar{z}^n=JW(z^{2n})$, $WJ(z^{2n-1})=0=JW(z^{2n-1})$, $W^*J(z^n)=\bar{z}^{2n}=JW^*(z^n)$, it follows that $K_\varphi=JA_\varphi=JWM_\varphi=WJM_\varphi=WS_\varphi$, where S_φ is a Hankel operator on L^2 .

- (b) $||K_{\varphi}|| = ||A_{\varphi}||$.
- (c) We note that J has a doubly infinite Hankel matrix, and an operator A having a Hankel matrix is characterized by the operator equation $V^*A = AV$, where V is the bilateral shift [11].
- (d) K_{φ}^* , the adjoint of $K_{\varphi}^{"}$, is given by $K_{\varphi}^* = A_{\varphi}^*J^* = M_{\bar{\varphi}}W^*J = M_{\bar{\varphi}}JW^* = JM_{\bar{\varphi}(\bar{z})}W^* = JA_{\varphi(\bar{z})}^*$.

We know that an operator A on L^2 is a slant Toeplitz operator if and only if $VA = AV^2$, where V is the bilateral shift [7, Proposition 3]. We present here a similar characterization of a slant Hankel operator.

Theorem 2.2. An operator A on L^2 is a slant Hankel operator if and only if $V^*A = AV^2$, where V is the bilateral shift.

PROOF. Suppose $A=K_{\varphi}$ is a slant Hankel operator. Then $V^*K_{\varphi}=V^*JA_{\varphi}=JVA_{\varphi}=JA_{\varphi}V^2=K_{\varphi}V^2$. Conversely, suppose $V^*A=AV^2$. Then $VJA=JV^*A=JAV^2$. Therefore, JA is a slant Toeplitz operator on L^2 by [7, Proposition 3]. Consequently, A is a slant Hankel operator on L^2 .

COROLLARY 2.3. The set of all slant Hankel operators on L^2 is a subspace of $\mathbf{B}(L^2)$.

PROOF. If a and b are complex numbers and $\varphi, \psi \in L^{\infty}$, then

$$aK_{\varphi} + bK_{\psi} = aJA_{\varphi} + bJA_{\psi} = J(aA_{\varphi}) + J(bA_{\psi}) = J(aA_{\varphi} + bA_{\psi})$$
$$= J(A_{a\varphi+b\psi}) = K_{a\varphi+b\psi}.$$

Therefore, it is a linear manifold.

Suppose that for each α , K_{α} is a slant Hankel operator such that $K_{\alpha} \to K$ weakly, where $\{\alpha\}$ is a net. Then, for f,g in L^2 , we have $\langle K_{\alpha}V^2f,g\rangle \to \langle KV^2f,g\rangle$ and $\langle V^*K_{\alpha}f,g\rangle = \langle K_{\alpha}f,Vg\rangle \to \langle Kf,Vg\rangle = \langle V^*Kf,g\rangle$. Since $K_{\alpha}V^2 = V^*K_{\alpha}$ for all α , we get $\langle KV^2f,g\rangle = \langle V^*Kf,g\rangle$. This implies that $V^*K = KV^2$ and hence K is a slant Hankel operator by Theorem 2.2. Therefore, the set of all slant Hankel operators is weakly closed and hence strongly closed [5, Problem 13]. This completes the proof.

3. Compressions of slant Hankel operators

We denote the compression of a slant Hankel operator K_{φ} to H^2 by L_{φ} . By the definition of compression, we have $L_{\varphi} = PK_{\varphi}|_{H^2}$, equivalently, $L_{\varphi}P = PK_{\varphi}P$, where P is the orthogonal projection of L^2 onto H^2 . We have the following.

Theorem 3.1. $L_{\varphi} = WH_{\varphi}$, where H_{φ} is a Hankel operator on H^2 . (Note that $H_{\varphi} = PS_{\varphi}|_{H^2}$)

PROOF.
$$L_{\varphi} = PK_{\varphi}|_{H^2} = PJA_{\varphi}|_{H^2} = PJWM_{\varphi}|_{H^2} = WPJM_{\varphi}|_{H^2} = WH_{\varphi}.$$

Remark 3.2. (a) If $\varphi - \psi$ is in zH^{∞} , then for f in H^2 , we have $L_{\varphi - \psi}(f) = WH_{\varphi - \psi}(f) = WPJ((\varphi - \psi)f) = 0$, since $J((\varphi - \psi)f) = (\varphi - \psi)f(\bar{z})$ is in $H^{2^{\perp}}$. Therefore, $L_{\varphi} = L_{\psi}$. This implies that the mapping $\varphi \to L_{\varphi}$ is not one-one and hence φ is not unique.

- (b) If $\varphi(z) = 1$, then, for f in H^2 , we have $L_1(f) = WH_1(f) = WPJ(f) = WP(f(\bar{z})) = \langle f, z^0 \rangle z^0$. Hence L_1 is the projection of H^2 onto the subspace spanned by z^0 .
- (c) For f in H^2 , we have, by Theorem 2 [12], $H_{\varphi}W(f) = PJM_{\varphi}W(f) = PJWM_{\varphi(z^2)}(f) = PJA_{\varphi(z^2)}(f) = PK_{\varphi(z^2)}(f)$. Therefore, $H_{\varphi}W = L_{\varphi(z^2)}$.
- Z. Nehari [8] proved that an operator B on H^2 is a Hankel operator on H^2 if and only if $U^*B = BU$, where U is the unilateral shift. We state and prove a similar result for the compression of a slant Hankel operator. To achieve this we

need the 'lifting theorem' of Sz-Nagy and Foias [3], [4], [9] and [11]. One version of the theorem is as follows:

LIFTING THEOREM. For i=1,2, let B_i be a contraction on a Hilbert space H_i , and let A_i , acting on the Hilbert space K_i , be the minimal unitary dilation of B_i . Let P_i be the orthogonal projection of K_i onto H_i . Then an operator X from H_1 to H_2 satisfies $B_2X = XB_1$ only if there exists an operator Y from K_1 to K_2 such that (i) $A_2Y = YA_1$, (ii) $\|X\| = \|Y\|$, (iii) $P_2YP_1 = XP_1$.

Theorem 3.3. An operator L on H^2 is the compression of a slant Hankel operator if and only if $U^*L = LU^2$, where U is the unilateral shift. In that case ||L|| = ||K||, where $L = PK|_{H^2}$.

PROOF. Since $P(\bar{z}Wf) = PW(\bar{z}^2f) = WP(\bar{z}^2f)$, for f in H^2 , we have $U^*W = WU^{*2}$. Now, suppose $L = L_{\varphi}$, the compression of a slant Hankel operator. Then $L_{\varphi} = WH_{\varphi}$ and $U^*L_{\varphi} = U^*WH_{\varphi} = WU^{*2}H_{\varphi} = WH_{\varphi}U^2 = L_{\varphi}U^2$. For the converse, we first note that V, the bilateral shift, is the minimal unitary

For the converse, we first note that V, the bilateral shift, is the minimal unitary dilation of U, the unilateral shift; and V^* is the minimal unitary dilation of U^* [5, Problem 155]. Suppose $U^*L = LU^2$. Then by the lifting theorem, there is an operator K on L^2 such that $V^*K = KV^2$, ||K|| = ||L|| and LP = PKP. By Theorem 2.2, we get $K = K_{\varphi}$, for some φ in L^{∞} . Therefore, $PK_{\varphi}P = L_{\varphi}P$. Consequently, $L = L_{\varphi}$, the compression of K_{φ} . This completes the proof.

We give another proof of Theorem 3.2 by using S. Parrott's observation [10] which is as follows.

PARROTT'S OBSERVATION. The smallest norm of an operator matrix $\begin{pmatrix} X & C \\ B & A \end{pmatrix}$, as X varies, is given as the maximum of the norms of $\begin{pmatrix} 0 & 0 \\ B & A \end{pmatrix}$, and $\begin{pmatrix} 0 & C \\ 0 & A \end{pmatrix}$. Now, suppose that L is an operator such that $U^*L = LU^2$. Then $L = (a_{-2i-j})_{i,j=0}^{\infty}$. Let

$$K_{2,1} = \begin{pmatrix} a_2 & a_1 & a_0 & a_{-1} & a_{-2} & \bullet \\ a_0 & a_{-1} & a_{-2} & a_{-3} & a_{-4} & \bullet \\ a_{-2} & a_{-3} & a_{-4} & a_{-5} & a_{-6} & \bullet \\ \bullet & \bullet & \bullet & \bullet & \bullet \end{pmatrix} = \begin{pmatrix} X & C \\ B & A \end{pmatrix}.$$

Then

$$\begin{pmatrix} 0 & 0 \\ B & A \end{pmatrix} = \begin{pmatrix} 0 & C \\ 0 & A \end{pmatrix} = K_{2,1}$$

Therefore, by Parrott's observation, we have $||K_{2,1}|| = ||L||$. Consequently $K_{2,1}$ is bounded. Continuing this construction, let

$$K_{4,3} = \begin{pmatrix} a_4 & a_3 & a_2 & a_1 & a_0 & \bullet \\ a_2 & a_1 & a_0 & a_{-1} & a_{-2} & \bullet \\ a_0 & a_{-1} & a_{-2} & a_{-3} & a_{-4} & \bullet \\ \bullet & \bullet & \bullet & \bullet & \bullet \end{pmatrix} = \begin{pmatrix} X & C \\ B & A \end{pmatrix}.$$

Then

$$\begin{pmatrix} 0 & 0 \\ B & A \end{pmatrix} = \begin{pmatrix} 0 & C \\ 0 & A \end{pmatrix} = K_{2,1}$$

Therefore, by Parrott's observation again, we have $||K_{4,3}|| = ||K_{2,1}|| = ||L||$. Consequently $K_{4,3}$ is bounded. Continuing this construction, at the n-th step, a bounded linear transformation $K_{2n,2n-1}$ is constructed with $||K_{2n,2n-1}|| = ||L||$. It follows that $K = (a_{-2i-j})_{i,j=-\infty}^{\infty}$ and K is an operator on L^2 . Moreover, $V^*K = KV^2$. Therefore, by Theorem 2.2, $K = K_{\varphi}$, for some φ in L^{∞} and $L = PK_{\varphi}|_{H^2}$. Therefore, $L = L_{\varphi}$. This completes the proof.

Remark 3.4. According to the construction above it is also apparent that φ is not unique. In fact, as remarked earlier, if $\varphi - \psi$ is in zH^{∞} , then $L = L_{\varphi} = L_{\psi}$.

We observe that $||A_{\varphi}|| = ||W|\varphi|^2||_{\infty}^{1/2}$ [7, Proposition 5] and $||K_{\varphi}|| = ||A_{\varphi}||$ by Remark 2.1. Hence $||K_{\varphi}|| = ||W|\varphi|^2||_{\infty}^{1/2}$.

Theorem 3.5. We have $||L_{\varphi}|| = \inf\{||W|\varphi - \phi|^2||_{\infty}^{1/2} : \phi \in zH^{\infty}\}.$

PROOF. By Theorem 3.2, we know that there is a φ in L^{∞} and $\varphi - \psi$ in zH^{∞} such that $||L_{\varphi}|| = ||K_{\psi}|| = ||W|\psi|^2|_{\infty}^{1/2}$. This implies that

$$\inf\{\|W|\varphi - \phi|^2\|_{\infty}^{1/2} : \phi \in zH^{\infty}\} \leqslant \|W|\psi|^2\|_{\infty}^{1/2} = \|K_{\psi}\| = \|L_{\varphi}\|.$$

On the other hand, $||L_{\varphi}|| = ||L_{\varphi-\phi}|| \le ||K_{\varphi-\phi}|| = ||W|\varphi-\phi|^2||_{\infty}^{1/2}$. This implies that $||L_{\varphi}|| \le \inf\{||W|\varphi-\phi|^2||_{\infty}^{1/2} : \phi \in zH^{\infty}\}$. This completes the proof. \square

For f,g in H^2 , we have $\langle H_{\varphi}^*f,g\rangle=\langle f,H_{\varphi}g\rangle=\langle f,PJ(\varphi g)\rangle=\langle \bar{\varphi}(z)f(\bar{z}),g\rangle.$ This implies that $H_{\varphi}^*f=P(\bar{\varphi}(z)f(\bar{z}))=PJ(\bar{\varphi}(\bar{z})f(z))=H_{\bar{\varphi}(\bar{z})}f.$ Therefore, $H_{\varphi}^*=H_{\bar{\varphi}(\bar{z})}.$ Since $L_{\varphi}=WH_{\varphi},$ we have $L_{\varphi}^*=H_{\bar{\varphi}(\bar{z})}W^*.$

Theorem 3.6. $0 \neq L_{\varphi}$ is self-adjoint if and only if $\varphi(z) = non\text{-}zero$ real constant.

PROOF. If $\sum_{i=-\infty}^{\infty} a_i z^i$ is the Fourier expansion of φ , then the (i,j)-th entry of the matrix of L_{φ} is given by $\langle L_{\varphi} z^j, z^i \rangle = \langle W H_{\varphi} z^j, z^i \rangle = \langle P J(\varphi z^j), z^{2i} \rangle = \langle \varphi, z^{-2i-j} \rangle = a_{-2i-j}$.

Now, suppose L_{φ} is self-adjoint. Then, for $i, j \geq 0$, $a_{-2i-j} = \bar{a}_{-2j-i}$. Put i = 0. Then we have $a_{-j} = \bar{a}_{-2j}$. This implies that for each $k \geq 0$ and for all $n \geq 0$ $|a_{-k}| = |a_{-k2n}|$. This in turn implies that $a_{-k} = 0$, for all k > 0, because $a_{-k} \to 0$, as $k \to \infty$. Therefore, $\varphi(z) = a_0$.

Conversely, if $\varphi(z) = a_0$, then $L_{\varphi}(f) = WPJ(\varphi f) = a_0 \langle f, z^0 \rangle z^0$ and $L_{\varphi}^*(f) = H_{\bar{\varphi}(\bar{z})}W^*(f) = PJ(\bar{\varphi}(\bar{z})f(z^2)) = \bar{a}_0 \langle f(z^2), z^0 \rangle z^0 = \bar{a}_0 \langle f, z^0 \rangle z^0$. Since $\bar{a}_0 = a_0$, we have the desired result.

Remark 3.7. If $\varphi \in zH^{\infty}$, then $L_{\varphi} = 0$. Therefore, L_{φ} is self-adjoint.

By making the same type of calculations as in the proof of Theorem 3.5, we can prove the following.

Theorem 3.8. L_{φ} is hyponormal if and only if φ is in H^{∞} .

PROOF. Suppose L_{φ} is hyponormal. Then $L_{\varphi} = WH_{\varphi}$ and for f in H^2 , $||WH_{\varphi}f|| \ge ||H_{\varphi}^*W^*f||$. Equivalently, $||WPJ(\varphi f)|| \ge ||PJ(\bar{\varphi}(\bar{z})f(z^2))||$. Putting f(z) = 1, we get $||WPJ(\varphi)||^2 \ge ||PJ(\bar{\varphi}(\bar{z}))||^2$. Equivalently,

$$\sum_{i=0}^{\infty} |a_{-2i}|^2 \geqslant \sum_{i=0}^{\infty} |\bar{a}_{-i}|^2,$$

where $\sum_{i=-\infty}^{\infty} a_i z^i$ is the Fourier expansion of φ . This implies that $a_{-2i-1}=0$, for $i=0,1,2,\ldots$ Again putting f(z)=z, we get $\|WPJ(\varphi z)\|^2 \geqslant \|PJ(z^2\bar{\varphi}(\bar{z}))\|^2$. Equivalently,

$$\sum_{i=0}^{\infty} |a_{-2i-1}|^2 \geqslant \sum_{i=0}^{\infty} |\bar{a}_{-i-2}|^2.$$

But the left-hand side is equal to 0. Therefore $a_{-i-2}=0$, for $i=0,1,2,\ldots$. Consequently, $a_{-i}=0$, for $i=1,2,3,\ldots$, which means φ is in H^{∞} .

Conversely, let φ be in H^{∞} . Then $L_{\varphi} = 0$ if $\varphi \in zH^{\infty}$, and L_{φ} is a multiple of the projection on the subspace of H^2 spanned by z^0 if $\varphi(z) = \text{constant}$. In other words, if $\varphi(z) = a_0$, then $L_{\varphi}(f) = L_{a_0}(f) = a_0 \langle f, z^0 \rangle z^0$ and its adjoint $L_{\varphi}^*(f) = L_{a_0}^*(f) = \bar{a}_0 \langle f, z^0 \rangle z^0$. Therefore, L_{φ} is normal and hence hyponormal. This completes the proof.

REMARK 3.9. (a) The non-zero hyponormal L_{φ} are the scalar multiples of the projection of H^2 onto the subspace spanned by z^0 .

(b) A hyponormal L_{φ} is necessarily normal.

Theorem 3.10. L_{φ} can not be an isometry.

PROOF. Suppose L_{φ} is an isometry. Then, for $j=0,\ 1,2,\ldots,$ we have $\|L_{\varphi}z^j\|=\|z^j\|=1$. Equivalently,

$$\sum_{k=0}^{\infty} |a_{-2k-j}|^2 = 1,$$

where $\sum_{k=-\infty}^{\infty} a_k z^k$ is the Fourier expansion of φ . Putting j=0 and j=2, we get

$$\sum_{k=0}^{\infty} |a_{-2k}|^2 = \sum_{k=0}^{\infty} |a_{-2k-2}|^2 = 1.$$

This implies that $a_0=0$. In general, by putting j=2n and j=2n+2, we get $a_{-2n}=0$, for $n=0,1,2,\ldots$. Similarly by putting j=2n+1 and j=2n+3, we get $a_{-2n-1}=0$, for $n=0,1,2,\ldots$. Therefore, $\varphi(z)=\sum_{k=1}^\infty a_k z^k$, but this φ induces the zero operator, that is, $L_{\varphi}=0$. This is a contradiction. Hence L_{φ} cannot be an isometry.

Theorem 3.11. L_{φ} is never a Fredholm operator.

PROOF. Suppose L_{φ} is a Fredholm operator. Then: (i) $\operatorname{ran}(L_{\varphi})$ is closed, (ii) $\operatorname{dim} \ker(L_{\varphi})$ and $\operatorname{dim} \ker(L_{\varphi}^*)$ are finite.

If $\ker(L_{\varphi}) = \ker(L_{\varphi}^*) = \{0\}$, then L_{φ} would be invertible, and hence $U^* = L_{\varphi}U^2L_{\varphi}^{-1}$, as $U^*L_{\varphi} = L_{\varphi}U^2$ by Theorem 3.2. But this is not true, because U^* is not similar to U^2 . Therefore, either $\ker(L_{\varphi}) \neq \{0\}$ or $\ker(L_{\varphi}^*) \neq \{0\}$. Suppose $\ker(L_{\varphi}) \neq \{0\}$. Then there is a non-zero f in H^2 such that $L_{\varphi}f = 0$. Since $U^{*n}L_{\varphi} = L_{\varphi}U^{2n}$, by repeated use of Theorem 3.2, it follows that $U^{2n}f$ is in $\ker(L_{\varphi})$, for all $n = 1, 2, 3, \ldots$ Since $U^{2n}f$ are linearly independent for different n's, we have dim $\ker(L_{\varphi})$ is equal to infinity, and hence L_{φ} is not Fredholm. Similarly, if $\ker(L_{\varphi}^*) \neq \{0\}$, then there is a non-zero g in H^2 such that $L_{\varphi}g = 0$. Since $L_{\varphi}^*U^n = U *^{2n}L_{\varphi}^*$ by Theorem 3.2, it follows that U^ng is in $\ker(L_{\varphi}^*)$ and $\dim \ker(L_{\varphi}^*) = \infty$. Therefore, L_{φ} is not Fredholm. This completes the proof.

Consider the matrix of L_{φ}^* , the adjoint of L_{φ} , given below

$$\begin{pmatrix} \bar{a}_{0} & \bar{a}_{-2} & \bar{a}_{-4} & \bullet \\ \bar{a}_{-1} & \bar{a}_{-3} & \bar{a}_{-5} & \bullet \\ \bar{a}_{-2} & \bar{a}_{-4} & \bar{a}_{-6} & \bullet \\ \bullet & \bullet & \bullet & \bullet \end{pmatrix}$$

Since W eliminates every odd row of the matrix of L_{φ}^* , it follows that the matrix of WL_{φ}^* is a matrix of a Hankel operator as shown below

$$\begin{pmatrix}
\bar{a}_0 & \bar{a}_{-2} & \bar{a}_{-4} & \bullet \\
\bar{a}_{-2} & \bar{a}_{-4} & \bar{a}_{-6} & \bullet \\
\bar{a}_{-4} & \bar{a}_{-6} & \bar{a}_{-8} & \bullet \\
\bullet & \bullet & \bullet
\end{pmatrix}$$

If $\sum_{i=-\infty}^{\infty} a_i z^i$ is the Fourier expansion of φ , then the matrix above defines the Hankel operator induced by the function $W(\bar{\varphi}(\bar{z}))$. Therefore, $WL_{\varphi}^* = H_{\psi}$, where $\psi = W(\bar{\varphi}(\bar{z}))$.

REMARK 3.12. (a) If L_{φ} is compact, $then L_{\varphi}^*$ is also compact. By the above relation $WL_{\varphi}^* = H_{\psi}$, and hence H_{ψ} is compact. By Hartman's theorem [2] and [6], we have that ψ belongs to $H^{\infty} + C(T)$.

(b) If φ is in $H^{\infty} + C(T)$, then L_{φ} is also compact, since $L_{\varphi} = WH_{\varphi}$.

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