# TWO NEW FINITE DIFFERENCE METHODS FOR COMPUTING EIGENVALUES OF A FOURTH ORDER LINEAR BOUNDARY VALUE PROBLEM 

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#### Abstract

This paper describes some new finite difference methods of order 2 and 4 for computing eigenvalues of a two-point boundary value problem associated with a fourth order differential equation of the form ( $\left.p y^{\prime \prime}\right)^{\prime \prime}+(q-\lambda r) y=0$. Numerical results for two typical eigenvalue problems are tabulated to demonstrate practical usefulness of our methods.


KEY WORDS AND PHRASES. Band-matrices, finite-difference methods, generalized eigenvalue problem, positive definite matrices, two-point boundary value problems.
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## 1. INTRODUCTION.

We shall consider the fourth order linear differential equation

$$
\begin{equation*}
\frac{d^{2}}{d x^{2}}\left[p(x) \frac{d^{2} y}{d x^{2}}\right]+[q(x)-\lambda r(x)] y=0,-\infty \leq a \leq x<b<\infty \tag{1.1}
\end{equation*}
$$

associated with the following pairs of homogeneous boundary conditions

$$
\begin{equation*}
y(a)=y(b)=y^{\prime \prime}(a)=y^{\prime \prime}(b)=0 \tag{1.2}
\end{equation*}
$$

Such boundary value problens occur in applied mathematics, engineering and modern physics, (see ref. [1-4]. In the differential equation (1.1) the functions $p(x), q(x)$, $r(x) \in C[a, b]$ and satisfy the conditions

$$
\begin{equation*}
p(x)>0, q(x) \geq 0 \text { and } r(x)>0, x \in[a, b] \tag{1.3}
\end{equation*}
$$

We cannot compute the exact values of the eigenvalues $\lambda$ for which the boundary value problem (1.1) - (1.2) has a nontrivial eigensolution $y(x)$ for arbitrary chocies of the functions $p(x), q(x)$ and $r(x)$. We resort to numerical methods for computing approximate values of $\lambda$. The most commonly used technique for approximating $\lambda$ for which the system (1.1) - (1.2) has a nontrivial eigenfunction $y(x)$ is by finite difference methods.

Recently, the author [2] has analysed some new finite different methods of order 2 and 4 for computing eigenvalues of a two point boundary value problem involving the differential equation (1.1) with $p(x) \equiv 1$ associated with one of the following pairs of homogeneous boundary conditions:
(a) $y(a)=y(b)=y^{\prime}(a)=y^{\prime}(b)=0$
(b) the same boundary conditions as (1.2)
(c) $y(a)=y^{\prime}(a)=y^{\prime \prime}(b)=y^{\prime \prime \prime}(b)=0$.

Chawla and Katti [3] have developed a numerical finite difference method of order 2 for approximating the lowest eigenvalue $\lambda$ of the system (1.1) - (1.4(a)) with $p(x)$ $\equiv 1$. A fourth order method was later developed by Chawla [4] for the numerical treatment of the same problem. This latter method leads to a generalized seven-band symmetric matrix eigenvalue problem.

Let $\lambda$ be any eigenvalue of the system (1.1) - (1.2) and let $y(x) \neq 0$ be the corresponding eigenfunction. Then on multiplying (1.1) by $y(x)$ and integrating the resulting equation from $a$ to $b$, we find after integration by parts and on using (1.2), that

$$
\begin{equation*}
\lambda=\frac{\int_{a}^{b} p\left(y^{\prime \prime}\right)^{2} d x+\int_{a}^{b} q y^{2} d x}{\int_{a}^{b} r y^{2} d x}>0 \tag{1.5}
\end{equation*}
$$

in view of (1.3).
The purpose of this brief report is to present two new finite difference methods for computing approximate values of $\lambda$ for the system (1.1) - (1.2). These methods lead to generalized five-band and nine-band symmetrix matrix eigenvalue problems and provide $O\left(h^{2}\right)$ and $O\left(h^{4}\right)$-convergent approximations for the eigenvalues.
2. A SECOND ORDER METHOD

For a positive integer $N \geq 5$, let $h=(b-a) /(N+1)$ and $x_{i}=a+i h$, $i=0(1) N+1$. We shall designate $y_{i}=y\left(x_{i}\right), p_{i}=p\left(x_{i}\right), q_{i}=q\left(x_{i}\right)$ and $r_{i}=r\left(x_{i}\right)$. Note that the differential system (1.1) - (1.2) is equivalent to
(a) $y^{\prime \prime}(x)=v(x) / p(x), y(a)=y(b)=0$,
(b) $v^{\prime \prime}(x)+[q(x)-\lambda r(x)] y(x)=0$,

$$
\begin{equation*}
v(a)=v(b)=0 \tag{2.1}
\end{equation*}
$$

Now the central difference approximation to 2.1(a) is

$$
\begin{align*}
& -y_{i-1}+2 y_{i}-y_{i+1}+h^{2}\left(v_{i} / p_{i}\right)+\frac{h^{4}}{12} y^{(4)}\left(\theta_{i}\right)=0  \tag{2.2}\\
& \theta_{i} \in\left(x_{i-1}, x_{i+1}\right), i=1(1) N
\end{align*}
$$

The preceding system can be conveniently written in matrix form

$$
\begin{equation*}
J Y+h^{2}{ }_{p}^{-1} V+\frac{h^{4}}{12} T_{1}=0 \tag{2.3}
\end{equation*}
$$

where $\mathrm{Y}=\left(\mathrm{y}_{\mathrm{i}}\right), \mathrm{V}=\left(\mathrm{v}_{\mathrm{i}}\right), \mathrm{T}_{1}=\left(\rho_{\mathrm{i}}\right)$ are N -dimensional column vectors with $\rho_{i}=y^{(4)}\left(\theta_{i}\right), P=\operatorname{diag}\left(p_{i}\right)$, and $J=\left(j_{m n}\right)$ is a tridiagonal matrix so that

$$
j_{m}^{n}=\left\{\begin{align*}
2, & m=n  \tag{2.4}\\
-1, & |m-n|=1 \\
0, & |m-n|>1 .
\end{align*}\right.
$$

In an analogous manner, on discretizing 2.1(b), we get

$$
\begin{equation*}
J V-h^{2} Q Y+\lambda h^{2} R Y+\frac{h^{4}}{12} T_{2}=0 \tag{2.5}
\end{equation*}
$$

where $Q=\operatorname{diag}\left(q_{i}\right), R=\operatorname{diag}\left(r_{i}\right)$ and $T_{2}=\left(\sigma_{i}\right)$ with $\sigma_{i}=v^{(4)}\left(\phi_{i}\right)$, $\phi_{i} \in\left(x_{i-1}, x_{i+1}\right)$. Next, we eliminate $v$ between (2.3) and (2.5) to obtain

$$
\begin{equation*}
A Y \equiv\left(J P J+h^{4} Q\right) Y=\lambda h^{4} R Y+\Gamma, \tag{2.6}
\end{equation*}
$$

where

$$
\begin{equation*}
r=\frac{1}{12}\left[h^{6} \mathrm{~T}_{2}-h^{4} \mathrm{JPT}_{1}\right] \tag{2.7}
\end{equation*}
$$

It can be verified that the matrix $A=J P J+h^{4} Q$ is a five-band symmetric matrix. Now, in (2.6), neglect truncation error $\Gamma$, replace $Y$ by $\tilde{Y}$, then our method for computing approximations $\wedge$ for $\lambda$ of the system (1.1) - (1.2) can be expressed as a generalized seven-band symmetric matrix eigenvalue problem

$$
\begin{equation*}
A \tilde{Y}=\Lambda h^{4} R \tilde{Y} \tag{2.8}
\end{equation*}
$$

In fact the matrix JPJ is a positive definite matrix and hence for any step-size $h>0$, the approximations $\wedge$ for $\lambda$ by (2.8) are real and positive for all $p(x)>0$ and $r(x)>0$. That our method provides $0\left(h^{2}\right)$ convergent approximations $\wedge$ for $\lambda$ can be established following Grigorieff [5]. We omit the proof of convergence for brevity.
3. A FOURTH ORDER METHOD

Following Shoosmith [6] the boundary value problems 2.1(a) and 2.1(b) are discretized by the finite difference scheme
(a) $14 y_{0}-29 y_{1}+16 y_{2}-y_{3}=h^{2}\left[y_{0}^{\prime \prime}+12 y_{1}^{\prime \prime}\right]$,
(b) (1- $\left.\frac{\delta^{2}}{12}\right) \delta^{2} y_{i}=h^{2} y_{i}^{\prime \prime} \quad, \quad i=2(1) N-1$,
(c) $-\mathrm{y}_{\mathrm{n}-2}+16 \mathrm{y}_{\mathrm{N}-1}-29 \mathrm{y}_{\mathrm{N}}+14 \mathrm{y}_{\mathrm{N}+1}=\mathrm{h}^{2}\left[12 \mathrm{y}_{\mathrm{N}}^{\prime \prime}+\mathrm{y}_{\mathrm{N}+1}^{\prime \prime}\right]$.

It turns out the boundary value problem 3.1(a) gives rise to the linear equations

$$
\begin{equation*}
M \tilde{Y}+12 h_{p}^{2}-1 \tilde{V}=0 \tag{3.2}
\end{equation*}
$$

Similarly, for the system 2.1(b), we obtain the linear equations

$$
\begin{equation*}
M \tilde{V}=12 h^{2} Q \tilde{Y}-12 \Lambda h^{2} R \bar{Y} \tag{3.3}
\end{equation*}
$$

where the five-band $N \times N$ matrix $M$ is given by

$$
M=\left[\begin{array}{rrrrrrr}
29 & -16 & 1 & & &  \tag{3.4}\\
-16 & 30 & -16 & 1 & & \\
1 & -16 & 30 & -16 & 1 & & \\
-\cdots & - & - & -\cdots & - & - & - \\
& & 1 & -16 & 30 & -16 & 1 \\
& & & 1 & -16 & 30 & -16 \\
& & & & 1 & -16 & 29
\end{array}\right]
$$

The elimination of $\tilde{\mathrm{v}}$ from (3.2) and (3.3) gives our method for computing $\Lambda$ for $\lambda$ of (1.1) - (1.2) in the form

$$
\begin{equation*}
\left(M P M+144 h^{4} Q\right) \tilde{Y}=144 \Lambda h^{4} R \tilde{Y} \tag{3.5}
\end{equation*}
$$

where the matrix MPM is a nine-band positive definite matrix and hence for any step-size $h>0$, the approximations $\Lambda$ for $\lambda$ by (3.5) are real and positive for all $p(x), r(x)>0$. As before, it can be proved from the results of Grigorieff [5] that our present method provided $0\left(h^{4}\right)$ convergent approximations $\Lambda$ for $\lambda$.
4. NUMERICAL RESULTS

In order to illustrate our methods of order 2 and 4 for the approximation of $\lambda$ satisfying (1.1) - (1.2), we consider the eigenvalue problems:

$$
\begin{align*}
& {\left[\left(1+x^{2}\right) y^{\prime \prime}\right]^{\prime \prime}+\left[\frac{1}{\left(1+x^{2}\right)}-\lambda(1+x)^{4}\right] y=0,}  \tag{4.1}\\
& y(0)=y(1)=y^{\prime \prime}(0)=y^{\prime \prime}(1)=0 .
\end{align*}
$$

The smallest eigenvalue $\lambda_{1}=22.754,058,480$, . .

$$
\begin{align*}
& {\left[e^{x} y^{\prime \prime}\right]^{\prime \prime}+[\sin x-\lambda \cos x] y=0}  \tag{4.2}\\
& y(0)=y(1)=y^{\prime \prime}(0)=y^{\prime \prime}(1)=0
\end{align*}
$$

The smallest eigenvalue of the system (4.2) is $\lambda_{1}=181.345,488,233$, . . . We list the approximations $\Lambda_{1}$ for $\lambda_{1}$ and the relative errors $\left|1-\frac{\lambda_{1}}{\Lambda_{1}}\right|$ for various values of the step-size $h$. It is readily verified that the relative errors (Table I) based on generalized eigenvalue problem (2.8) provide $0\left(h^{2}\right)$ - convergent approximations for the smallest eigenvalue of the system (4.1) and (4.2). Similarly, the relative errors (Table II) based on the generalized eigenvalue problem (3.5) do indeed provide $0\left(h^{4}\right)$-convergent approximations for the smallest eigenvalue of the systems (4.1) and (4.2).

TABLE I

| Problem | N | $\Lambda_{1}$ | $11-\frac{\lambda_{1}}{\Lambda_{1}}$ |
| :---: | :---: | :---: | :---: |
| (4.1) | 7 | 22.187 | 2.557-2* |
|  | 15 | 22.610 | 6.352-3 |
|  | 31 | 22.718 | 1.586-3 |
|  | 63 | 22.745 | 3.962-4 |
|  | 127 | 22.752 | 9.907-5 |
|  | 255 | 22.753 | 2.480-5 |
| (4.2) | 7 | 176.641 | 2.664-2 |
|  | 15 | 180.159 | 6.588-3 |
|  | 31 | 181.048 | 1.642-3 |
|  | 63 | 181.271 | 4.103-4 |
|  | 127 | 181.327 | 1.025-4 |
|  | 255 | 181.341 | 2.560-5 |
| *We write 2.557-2 for $2.557 \times 10^{-2}$. |  |  |  |
| TABLE II |  |  |  |
| Results based on (3.5), 4th order approximations |  |  |  |
| Problem | N | $\Lambda_{1}$ | $\left\|1-\frac{\lambda^{1}}{\Lambda_{1}}\right\|$ |
| (4.1) | 7 | 22.746, 419 | 3.358-4 |
|  | 15 | 22.753, 574 | 2.129-5 |
|  | 31 | 22.754, 027 | 1.358-6 |
|  | 63 | 22.754, 056 | 1.078-7 |
| (4.2) |  |  |  |
|  | 15 | 181.339, 089 | 3.529-5 |
|  | 31 | 181.345, 093 | 2.175-6 |
|  | 63 | 181.345, 470 | 9.728-8 |

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