

Edgeworth expansions for a sample sum from a finite set of independent random variables*

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Abstract

Let $\{X_1, \dots, X_N\}$ be a set of N independent random variables, and let S_n be a sum of n random variables chosen without replacement from the set $\{X_1, \dots, X_N\}$ with equal probabilities. In this paper we give a one-term Edgeworth expansion of the remainder term for the normal approximation of S_n under mild conditions.

Key words: Edgeworth expansion, finite population, sampling without replacement.

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1 Introduction and main results

Let $\{X_1, \dots, X_N\}$ be a set of independent random variables, $\mu_k = EX_k$, $1 \leq k \leq N$. Let $R = (R_1, \dots, R_N)$ be a random vector independent of X_1, \dots, X_N , such that $P(R = r) = 1/N!$ for any permutation $r = (r_1, \dots, r_N)$ of the numbers $1, \dots, N$, and put $S_n = \sum_{j=1}^n X_{R_j}$, $1 \leq n \leq N$, that is for a sum of n random variables chosen without replacement from the set $\{X_1, \dots, X_N\}$ with equal probabilities.

In the situation that $X_k = \mu_k$, $1 \leq k \leq N$, are (nonrandom) real numbers, the sample sum S_n has been studied by a number of authors. The asymptotic normality was established by Erdős and Rényi (1959) under quite general conditions. The rate in the Erdős and Rényi central limit theorem was studied by Bikelis (1969) and later Höglund (1978). An Edgeworth expansion was obtained by Robinson (1978), Bickel and van Zwet (1978), Schnell (1989), Babu and Bai (1996) and later Bloznelis (2000a, b). Extensions to U -statistics and, more generally, symmetric statistics can be found in Nandi and Sen (1963), Zhao and Chen (1987, 1990), Kocic and Weber (1990), Bloznelis and Götze (2000, 2001) and Bloznelis (2003).

In contrast to rich investigations for the case $X_k = \mu_k$, $1 \leq k \leq N$, are (nonrandom) real numbers, there are only a few results concerned with the asymptotics of general S_n discussed in this paper. von Bahr (1972) showed that the distribution of $S_n/\sqrt{Var S_n}$ may be approximated by a normal distribution under certain mild conditions. The rate of the normal approximation has currently been established by Zhao, Wu and Wang (2004), in which the paper improved essentially earlier work by von Bahr (1972). Along the lines of Zhao, Wu and Wang (2004), this paper discusses Edgeworth expansions for the distribution of $S_n/\sqrt{Var S_n}$. Throughout the paper, let

$$\gamma_{12} = \frac{1}{N} \sum_{k=1}^N EX_k EX_k^2, \quad \alpha_j = \frac{1}{N} \sum_{k=1}^N (EX_k)^j, \quad \beta_j = \frac{1}{N} \sum_{k=1}^N E(X_k^j),$$

for $j = 1, 2, 3, 4$, and

$$p = n/N, \quad q = 1 - p, \quad b = 1 - p\alpha_2.$$

Theorem 1. *Suppose that $\alpha_1 = 0$ and $\beta_2 = 1$. Then, for all $1 \leq n < N$,*

$$\begin{aligned} & \sup_x \left| P(S_n/\sqrt{nb} \leq x) - G_n(x) \right| \\ & \leq C \left(\Delta_{1n} + (nq)^{-1} \right) + 3\sqrt{nq} \log(nq) \exp\{-nq\delta_N\}, \end{aligned} \quad (1)$$

where C is an absolute constant,

$$G_n(x) = \Phi(x) - \frac{\beta_3 - 3p\gamma_{12} + 2p^2\alpha_3}{6\sqrt{nb}^{3/2}} \Phi'''(x),$$

with $\Phi(x)$ being a standard normal distribution,

$$\Delta_{1n} = (nb)^{-1}\alpha_4 + \frac{(nb^2)^{-1}}{N} \sum_{k=1}^N E(X_k - pEX_k)^4,$$

and

$$\delta_N = 1 - \sup_{\delta_0 b / (9\mathcal{L}_0) \leq |t| \leq 16\sqrt{nb}} \left| \frac{1}{N} \sum_{k=1}^N E e^{itX_k} \right|,$$

where $\mathcal{L}_0 = \frac{1}{N} \sum_k E|X_k|^3$ and δ_0 is so small that $192\delta_0^2 + 24\delta_0 \leq 1 - \cos(1/16)$.

Property (1) improves essentially a result of Mirakhmedov (1983). The related result in Theorem 1 of Mirakhmedov (1983) depends on $\max_{1 \leq k \leq N} EX_k^4$. Note that it is frequently the case that $N^{-1} \sum_{k=1}^N EX_k^4$ is bounded, but $\max_{1 \leq k \leq N} EX_k^4$ tends to ∞ . Also note that Corollary 1 of Mirakhmedov (1983) requires $\overline{\lim}_{t \rightarrow \infty} |E e^{itX_k}| \leq \epsilon < 1$. This condition is quite restrictive since it takes away the most interesting case that the X_k are all degenerate.

When $X_k = \mu_k$, $1 \leq k \leq N$, are (nonrandom) real numbers, it is readily seen that $\alpha_2 = 1$, $b = q$, $\alpha_3 = \beta_3 = \gamma_{12} = \frac{1}{N} \sum_{k=1}^N \mu_k^3$,

$$\Delta_{1N} + (nq)^{-1} \leq 3(nq)^{-1} \frac{1}{N} \sum_{k=1}^N \mu_k^4.$$

In this case, the property (1) reduces to

$$\begin{aligned} & \sup_x |P(S_n/\sqrt{nq} \leq x) - G_{1n}(x)| \\ & \leq C(nq)^{-1} \frac{1}{N} \sum_{k=1}^N \mu_k^4 + 3\sqrt{nq} \log(nq) \exp\{-nq\delta_N\}, \end{aligned}$$

where $G_{1n} = \Phi(x) + \frac{p-q}{6\sqrt{nq}} \frac{1}{N} \sum_{k=1}^N \mu_k^3 \Phi'''(x)$, which gives one of main results in Bloznelis (2000a, b).

We next give a result complementary to Theorem 1. The result is better than Theorem 1 under certain conditions such as some of the X_k 's are non-degenerate random variables and q is close to 0.

Theorem 2. *Suppose that $\alpha_1 = 0$ and $\beta_2 = 1$. Then, for all $1 \leq n \leq N$,*

$$\sup_x \left| P(S_n/\sqrt{nb} \leq x) - G_n(x) \right| \leq C \Delta_{2n} + 3\sqrt{n} \log n \exp\{-n\delta_{1N}\}, \quad (2)$$

where C is an absolute constant, $G_n(x)$, \mathcal{L}_0 and δ_0 are defined as in Theorem 1, $\Delta_{2n} = (nb^2)^{-1} \beta_4$ and

$$\delta_{1N} = 1 - \sup_{\delta_0 b / (9\mathcal{L}_0) \leq |t| \leq 16\sqrt{nb}} \frac{1}{N} \sum_{k=1}^N |E e^{itX_k}|.$$

In the next section, we prove the main results. Throughout the paper we shall use C, C_1, C_2, \dots to denote absolute constants whose value may differ at each occurrence. Also, $I(A)$ denotes the indicator function of a set A , $\sharp(A)$ denotes the number of elements in the set A , \sum_k denotes $\sum_{k=1}^N$, and \prod_k denotes $\prod_{k=1}^N$. The symbol i will be used exclusively for $\sqrt{-1}$.

2 Proofs of Theorems

Let $\mu_k = EX_k$ and $\Psi(t) = E \exp\{itS_n/\sqrt{nb}\}$. Recall $\alpha_1 = 0$ and $\beta_2 = 1$. As in (4) of Zhao, Wu and Wang (2004),

$$\Psi(t) = [B_n(p)]^{-1} \int_{|\psi| \leq \pi\sqrt{nq}} \prod_k E\rho_k(\psi, t) d\psi, \quad (3)$$

where $B_n(p) = \sqrt{2\pi nq}G_n(p)$, $G_n(p) = \sqrt{2\pi}C_N^n p^n q^{N-n}$, $X_k^* = X_k - p\mu_k$ and

$$\rho_k(\psi, t) = q \exp\left\{-\frac{ip\psi}{\sqrt{nq}} - \frac{ip\mu_k t}{\sqrt{nb}}\right\} + p \exp\left\{\frac{iq\psi}{\sqrt{nq}} + \frac{itX_k^*}{\sqrt{nb}}\right\}.$$

The main idea of the proofs is outlined as follows. We first provide the expansions and the basic properties for $\prod_k E\rho_k(\psi, t)$ in Lemmas 1–4. In Lemma 5, the idea in von Bahr (1972) is extended to give an expansion of $\Psi(t)$ for the case $n/N \geq 1/2$. The proofs of Theorems 1 and 2 are finally completed by virtue of the classical Esseen's smoothing lemma.

In the proofs of Lemmas 1–4, we assume that $\Delta_{1n} < 1/16$ and $nq > 256$, where Δ_{1n} is defined as in Theorem 1. Throughout this section, we also define,

$$h(\psi, t) = \prod_k E\rho_k(\psi, t) \quad \text{and} \quad g(\psi, t) = \left(1 + \frac{i^3 f(\psi, t)}{6\sqrt{N}}\right) e^{-(\psi^2 + t^2)/2},$$

where $f(\psi, t) = A_0\psi^3 + 3A_1\psi t^2 + A_2t^3$, with

$$A_0 = \frac{q-p}{\sqrt{pq}}, \quad A_1 = \frac{(1-2p\alpha_2)\sqrt{pq}}{pb}, \quad A_2 = \frac{\beta_3 - 3p\gamma_{12} + 2p^2\alpha_3}{p^{1/2}b^{3/2}}. \quad (4)$$

Lemma 1. For $|\psi| \leq (nq)^{1/4}/4$ and $|t| \leq \Delta_{1n}^{-1/4}/4$, we have

$$|h(\psi, t) - g(\psi, t)| \leq C[\Delta_{1n} + (nq)^{-1}](s^4 + s^8) \exp\{-s^2/3\}, \quad (5)$$

where $s^2 = \psi^2 + t^2$.

Proof. Define a sequence of independent random vectors (U_k, V_k) , $1 \leq k \leq N$, by the conditional distribution given X_k^* as follows:

$$\begin{aligned} P(U_k = -p/\sqrt{pq}, V_k = -p\mu_k/\sqrt{pb} | X_k^*) &= q, \\ P(U_k = q/\sqrt{pq}, V_k = X_k^*/\sqrt{pb} | X_k^*) &= p. \end{aligned}$$

Let $W_k = \psi U_k + tV_k$. As in Lemma 1 of Zhao, Wu and Wang (2004), tedious but simple calculations show that

$$\begin{aligned} EW_k &= 0, \quad \sum_k EW_k^2 = N(\psi^2 + t^2), \quad \sum_k EW_k^3 = Nf(\psi, t), \\ \sum_k EW_k^4 &\leq 8 \sum_k E(\psi^4 U_k^4 + t^4 V_k^4) \leq 8N^2(\psi^4/(nq) + t^4 \Delta_{1n}). \end{aligned} \quad (6)$$

Furthermore, if we let $B_n^2 = \sum_k EW_k^2$, $\mathcal{L}_{jN} = \sum_k E|W_k|^j/B_n^j$, $j = 3, 4$, then

$$\mathcal{L}_{3N}^2 = \left(\sum_k E|W_k|^3 \right)^2 / B_n^6 \leq \mathcal{L}_{4N} \leq 8((nq)^{-1} + \Delta_{1n}), \quad (7)$$

and whenever $|\psi| \leq (nq)^{1/4}/4$ and $|t| \leq \Delta_{1n}^{-1/4}/4$,

$$s := \sqrt{\psi^2 + t^2} \leq \mathcal{L}_{4N}^{-1/4} [8(\psi^4/(nq) + t^4\Delta_{1n})]^{1/4} \leq \mathcal{L}_{4N}^{-1/4}/2. \quad (8)$$

Now, by recalling that W_k are independent r.v.s and noting that

$$h(\psi, t) = Ee^{i\sum_k W_k/\sqrt{N}} = Ee^{is\sum_k W_k/B_n},$$

it follows from (7)-(8) and the classical result (see, for example, Theorem 8.6 in Bhattacharya and Ranga Rao (1976)) that, for $|\psi| \leq (nq)^{1/4}/4$ and $|t| \leq \Delta_{1n}^{-1/4}/4$,

$$\begin{aligned} & \left| h(\psi, t) - \left(1 + \frac{i^3}{6\sqrt{N}} f(\psi, t) \right) e^{-s^2/2} \right| \\ &= \left| h(\psi, t) - \left(1 + \frac{i^3 s^3}{6B_n^3} \sum_k EW_k^3 \right) e^{-s^2/2} \right| \\ &\leq C(\mathcal{L}_{4N} + \mathcal{L}_{3N}^2)(s^4 + s^8)e^{-s^2/3} \\ &\leq C[\Delta_{1n} + (nq)^{-1}](s^4 + s^8) \exp\{-s^2/3\}. \end{aligned}$$

This proves (5) and hence completes the proof of Lemma 1. \square

Lemma 2. For $|\psi| \leq (nq)^{1/4}/4$ and $|t| \leq 1/4$, we have

$$\left| \frac{dh(\psi, t)}{dt} - \frac{dg(\psi, t)}{dt} \right| \leq C(1 + \psi^{12}) ((nq)^{-1} + \Delta_{1n}) e^{-\psi^2/4}. \quad (9)$$

Proof. We first show that if $|\psi| \leq (nq)^{1/4}/4$ and $|t| \leq 1/4$, then

$$\begin{aligned} \Lambda(\psi, t) &:= \left| \frac{dh(\psi, t)}{dt} + \left(t + \frac{i}{6\sqrt{N}} \frac{df(\psi, t)}{dt} \right) h(\psi, t) \right| \\ &\leq C(1 + \psi^4) ((nq)^{-1} + \Delta_{1n}) |h(\psi, t)|. \end{aligned} \quad (10)$$

To prove (10), define (U_k, V_k) and $W_k = \psi U_k + tV_k$ as in Lemma 1. Recall that $h(\psi, t) = E \exp\{i \sum_k W_k/\sqrt{N}\}$. It is readily seen that

$$\frac{dh(\psi, t)}{dt} = \frac{i h(\psi, t)}{\sqrt{N}} \sum_k I_{Nk}, \quad (11)$$

where $I_{Nk} = \left[E \exp\{iW_k/\sqrt{N}\} \right]^{-1} E \left[V_k \exp\{iW_k/\sqrt{N}\} \right]$. Recall $nq > 256$. It follows from (19) and (20) in Zhao, Wu and Wang (2004) that for $|\psi| \leq (nq)^{1/4}/4$ and $|t| \leq 1/4$,

$$\left[E \exp\{iW_k/\sqrt{N}\} \right]^{-1} = 1 + \theta_1 N^{-1} (\psi^2 + t^2 EV_k^2), \quad (12)$$

where $|\theta_1| \leq 1$ and $N^{-1}(\psi^2 + t^2 EV_k^2) \leq 1/4$. This, together with Taylor's expansion of e^{ix} , yields that (recall $EV_k = 0$)

$$\begin{aligned} & \left| I_{Nk} - \frac{i}{\sqrt{N}} EV_k W_k + \frac{1}{2N} EV_k W_k^2 \right| \\ & \leq \frac{3}{N^{3/2}} E|V_k| |W_k|^3 + \frac{\psi^2 + t^2 EV_k^2}{N} \left(\frac{1}{\sqrt{N}} |EV_k W_k| + \frac{1}{2N} |EV_k W_k^2| \right). \end{aligned} \quad (13)$$

As in the proof of (6), for $|t| \leq 1/4$,

$$\begin{aligned} \sum_k E|V_k| |W_k|^3 & \leq C(1 + |\psi|^3) \sum_k E(U_k^4 + V_k^4) \\ & \leq C(1 + |\psi|^3) N^2 [(nq)^{-1} + \Delta_{1n}], \end{aligned} \quad (14)$$

$$\begin{aligned} \sum_k (\psi^2 + EV_k^2) |EV_k W_k| & \leq C(1 + |\psi|^3) \sum_k (1 + EV_k^2) (EU_k^2 + EV_k^2) \\ & \leq C(1 + |\psi|^3) \sum_k (1 + EV_k^4) \\ & \leq C(1 + |\psi|^3) N^2 [(nq)^{-1} + \Delta_{1n}], \end{aligned} \quad (15)$$

$$\begin{aligned} \sum_k (\psi^2 + EV_k^2) |EV_k W_k^2| & \leq C(1 + \psi^4) \sum_k (1 + EV_k^2) (E|V_k|^3 + E|U_k|^3) \\ & \leq C(1 + \psi^4) \sum_k \left((pq)^{-1/2} (1 + EV_k^2) + E|V_k|^3 + EV_k^2 E|V_k|^3 \right) \\ & \leq C(1 + \psi^4) \left[N(pq)^{-1/2} + \sum_k \left(1 + EV_k^4 + \sqrt{N} EV_k^4 \right) \right] \\ & \leq C(1 + \psi^4) N^{5/2} [(nq)^{-1} + \Delta_{1n}], \end{aligned} \quad (16)$$

where, in the proof of (16), we have used the estimates: $|V_k|^3 \leq 1 + V_k^4$ and

$$EV_k^2 E|V_k|^3 \leq (EV_k^2)^{1/2} EV_k^4 \leq \sqrt{N} EV_k^4.$$

Now (10) follows from (11), (13)-(16) and

$$\sum_k EV_k W_k = tN, \quad \sum_k EV_k W_k^2 = N(2A_1 \psi t + A_2 t^2) = \frac{N}{3} \frac{df(\psi, t)}{dt}. \quad (17)$$

We next complete the proof of Lemma 2 by virtue of (10) and Lemma 1. We first notice that, by (6), for all ψ and t ,

$$\begin{aligned} |f(\psi, t)| & \leq \frac{1}{N} \sum_k E|W_k|^3 \leq \frac{1}{N} \left(\sum_k EW_k^2 \sum_k EW_k^4 \right)^{1/2} \\ & \leq 3\sqrt{N} (\psi^2 + t^2)^{3/2} (\Delta_{1n} + (nq)^{-1})^{1/2}, \end{aligned} \quad (18)$$

and similarly by (17), for all ψ and t ,

$$\left| \frac{df(\psi, t)}{dt} \right| \leq \frac{3}{N} \sum_k E|V_k W_k^2| \leq 9\sqrt{N} (\psi^2 + t^2)^{3/2} (\Delta_{1n} + (nq)^{-1})^{1/2}. \quad (19)$$

It follows from (18) and Lemma 1 that for $|\psi| \leq (nq)^{1/4}/4$ and $|t| \leq 1/4$,

$$|h(\psi, t) - e^{-(\psi^2+t^2)/2}| \leq C(\Delta_{1n} + (nq)^{-1})^{1/2} e^{-\psi^2/4}. \quad (20)$$

Therefore, by noting

$$\frac{dg(\psi, t)}{dt} = -tg(\psi, t) - \frac{i}{6\sqrt{N}} \frac{df(\psi, t)}{dt} e^{-(\psi^2+t^2)/2},$$

simple calculations show that

$$\begin{aligned} \left| \frac{dh(\psi, t)}{dt} - \frac{dg(\psi, t)}{dt} \right| &\leq \Lambda(\psi, t) + t|h(\psi, t) - g(\psi, t)| \\ &\quad + \frac{1}{6\sqrt{N}} \left| \frac{df(\psi, t)}{dt} \right| |h(\psi, t) - e^{-(\psi^2+t^2)/2}| \\ &\leq C(1 + \psi^{12})((nq)^{-1} + \Delta_{1n})e^{-\psi^2/4}, \end{aligned}$$

where we have used (5), (10), (19) and (20). The proof of Lemma 2 is now complete. \square

Lemma 3. Assume that $|\psi| \leq (nq)^{1/4}/4$. Then,

$$|h(\psi, t)| \leq C\Delta_{1n}e^{-(\psi^2+t^2)/4}, \quad (21)$$

for $\Delta_{1n}^{-1/4}/4 \leq |t| \leq (\Delta^*)^{-1}/16$, where

$$\Delta^* = N^{-1} \sum_k |\mu_k|^3 / \sqrt{nb} + N^{-1} \sum_k E|X_k - p\mu_k|^3 / (\sqrt{nb}^{3/2}).$$

Assume that $(nq)^{1/4}/4 \leq |\psi| \leq \pi\sqrt{nq}$. Then,

$$|h(\psi, t)| \leq C(nq)^{-4}, \quad (22)$$

for all $|t| \leq \delta_0(\Delta^*)^{-1}$, where δ_0 is so small that $192\delta_0^2 + 24\delta_0 \leq 1 - \cos(1/16)$. If in addition $|t| \leq 1/4$, then we also have

$$\left| \frac{dh(\psi, t)}{dt} \right| \leq C(nq)^{-4}. \quad (23)$$

Proof. The proof of this lemma follows directly from an application of Lemmas 1-3 in Zhao, Wu and Wang(2004). The choice of δ_0 can be found in the proof of Lemma 2 in Zhao, Wu and Wang(2004). We omit the details. \square

Lemma 4. Assume that $n/N \geq 1/2$ and $\Delta_{2n} \leq (nq)^{-1}/25$, where Δ_{2n} is defined as in Theorem 2. Then, for $|t| \leq \frac{1}{15}\sqrt{nb}^{3/2}/\mathcal{L}_0$,

$$|h(\psi, t)| \leq \exp\{-Ct^2\}, \quad (24)$$

where $\mathcal{L}_0 = \frac{1}{N} \sum_k E|X_k|^3$ is defined as in Theorem 1.

Proof. We only need to note that the condition $\Delta_{2n} \leq (nq)^{-1}/25$ implies that

$$5q\alpha_2 \leq 5q\beta_4^{1/2} \leq b = \frac{1}{N} \sum_k \text{Var}(X_k) + q\alpha_2,$$

that is, $\frac{1}{N} \sum_k \text{Var}(X_k) \geq (4/5)b$. Then (24) is obtained by repeating the proof of Lemma 4 in Zhao, Wu and Wang(2004). \square

Lemma 5. *Assume that $n/N \geq 1/2$ and $\Delta_{2n} \leq 1$. Then, for $|u| \leq \frac{1}{16} \min \{(n/\beta_4)^{1/4}, \frac{1}{8}b\sqrt{n}/\mathcal{L}_0\}$,*

$$\begin{aligned} & \left| E \exp\{iuS_n/\sqrt{n}\} - \exp\{-bu^2/2\} \left(1 + \frac{i^3 u^3 b^{3/2}}{6\sqrt{N}} A_2 \right) \right| \\ & \leq Cn^{-1} \beta_4(u^2 + u^4 + u^6 b) \exp\{-0.3bu^2\}, \end{aligned} \quad (25)$$

where $\mathcal{L}_0 = \frac{1}{N} \sum_k E|X_k|^3$ is defined as in Theorem 1.

Proof. Write, for $1 \leq k \leq N$,

$$f_k(u) = E \exp\{iuX_k/\sqrt{n}\}, \quad b_k(u) = \exp\{u^2/(2n)\}f_k(u) - 1$$

and $B_j = ((-1)^{j+1}/j) \sum_{k=1}^N b_k^j(u)$ for $1 \leq j \leq n$. As in von Bahr (1972), we have

$$\exp\{u^2/2\} E \exp\{iuS_n/\sqrt{n}\} = \sum_{i_j \geq 0, 1 \leq j \leq n} \prod_{j=1}^n \frac{(p^j B_j)^{i_j}}{i_j!} C_{N,n,\sum_{j=1}^n i_j}, \quad (26)$$

where

$$C_{N,n,r} = \begin{cases} \binom{N-r}{n-r} / \binom{N}{n}, & r \leq n; \\ 0, & r > n. \end{cases}$$

In view of (28) of Zhao, Wu and Wang (2004), for $r > 0$, $C_{N,n,r} \leq 1$, and for $n \geq 4$ and $r \leq n$,

$$C_{N,n,r} \geq 1 - r^2/n. \quad (27)$$

To prove (25) by using (26), we need some preliminary results.

Write $\beta_{jk} = EX_k^j, j = 2, 3, 4$. Recall that $N^{-1} \sum_k \beta_{2k} = 1$. We have that $\beta_4 \geq 1$ and by Taylor's expansion, for $|u| \leq \frac{1}{16}(n/\beta_4)^{1/4}$,

$$\begin{aligned} \exp\{u^2/(2n)\} &= 1 + \frac{1}{2n}u^2 + \frac{1}{8n^2}u^4 + \frac{\theta_4}{n^3}u^6, \quad \text{where } |\theta_4| \leq 1/24, \\ f_k(u) &= 1 + \frac{iu}{\sqrt{n}}\mu_k - \frac{u^2}{2n}\beta_{2k} - \frac{iu^3}{6n^{3/2}}\beta_{3k} + \frac{\theta_5 u^4}{n^2}\beta_{4k}, \quad \text{where } |\theta_5| \leq 1/24. \end{aligned}$$

Now, by noting that $|\mu_k| \leq \beta_{4k}^{1/4} \leq 1 + \beta_{4k}$, $|\beta_{2k}| \leq \beta_{4k}^{1/2} \leq 1 + \beta_{4k}$ and $|\beta_{3k}| \leq \beta_{4k}^{3/4} \leq 1 + \beta_{4k}$, we obtain that, for $|u| \leq \frac{1}{16}(n/\beta_4)^{1/4}$,

$$\begin{aligned} b_k(u) &= \exp\{u^2/(2n)\}f_k(u) - 1 \\ &= \frac{iu}{\sqrt{n}}\mu_k + \frac{u^2}{2n}(1 - \beta_{2k}) + \frac{iu^3}{6n^{3/2}}(3\mu_k - \beta_{3k}) + R_{1k}(u). \end{aligned} \quad (28)$$

where $|R_{1k}(u)| \leq (1 + \beta_{4k})u^4/n^2$. Furthermore, by noting that

$$\beta_{4k}^{1/4}|u|/\sqrt{n} \leq (N\beta_4)^{1/4}|u|/\sqrt{n} \leq 1/8$$

since $n/N \geq 1/2$ and $|u| \leq \frac{1}{16}(n/\beta_4)^{1/4}$, we have

$$b_k^2(u) = -\frac{u^2}{n}\mu_k^2 + \frac{iu^3}{n^{3/2}}\mu_k(1 - \beta_{2k}) + R_{2k}(u), \quad (29)$$

$$b_k^3(u) = -\frac{iu^3}{n^{3/2}}\mu_k^3 + R_{3k}(u), \quad (30)$$

$$|b_k^j(u)| \leq 3(1/2)^{j-4}(1 + \beta_{4k})u^4/n^2, \quad \text{for } j \geq 4, \quad (31)$$

where $|R_{2k}(u)| \leq 3(1 + \beta_{4k})u^4/n^2$ and $|R_{3k}(u)| \leq 4(1 + \beta_{4k})u^4/n^2$. Recalling that $\sum_k \mu_k = \sum_k (1 - \beta_{2k}) = 0$, it follows from (28)–(31) that, for $|u| \leq \frac{1}{16}(n/\beta_4)^{1/4}$,

$$pB_1 = p \sum_k b_k(u) = -\frac{iu^3}{6\sqrt{n}}\beta_3 + \theta_6\beta_4u^4/n, \quad (32)$$

$$p^2B_2 = (-p^2/2) \sum_k b_k^2(u) = \frac{u^2p}{2}\alpha_2 + \frac{iu^3p}{2\sqrt{n}}\gamma_{12} + \theta_7\beta_4u^4/n, \quad (33)$$

$$p^3B_3 = (p^3/3) \sum_k b_k^3(u) = -\frac{ip^2u^3}{3\sqrt{n}}\alpha_3 + \theta_8\beta_4u^4/n, \quad (34)$$

$$|p^jB_j| \leq \sum_k |b_k^j(u)| \leq 6\beta_4u^4(1/2)^{j-4}/n, \quad \text{for } j \geq 4. \quad (35)$$

where $|\theta_6| \leq 2$, $|\theta_7| \leq 3$ and $|\theta_8| \leq 3$. By virtue of (35), it is readily seen that, for $|u| \leq \frac{1}{16}(n/\beta_4)^{1/4}$,

$$\sum_{j=4}^n |p^jB_j| \leq 12\beta_4u^4/n, \quad (36)$$

Noting that $|\alpha_3| + |\beta_3| + |\gamma_{12}| \leq 3\mathcal{L}_0$ and recalling that $\Delta_{2n} = (nb^2)^{-1}\beta_4 \leq 1$, it follows easily from (32)–(34) and (36) that, for $|u| \leq \frac{1}{16} \min \{(n/\beta_4)^{1/4}, \frac{1}{8}b\sqrt{n}/\mathcal{L}_0\}$,

$$\begin{aligned} \sum_{j=1}^n |p^jB_j| &= \frac{1}{2}p\alpha_2u^2 + \theta_9\frac{u^3\mathcal{L}_0}{\sqrt{n}} + \theta_{10}\frac{u^4\beta_4}{n} \\ &= \frac{1}{2}p\alpha_2u^2 + \theta_{11}bu^2, \end{aligned} \quad (37)$$

where $|\theta_9| \leq 3$, $|\theta_{10}| \leq 20$ and $|\theta_{11}| \leq 0.2$. Also, if we let $L(u) = \sum_{j=1}^3 p^jB_j - p\alpha_2u^2/2$, we have

$$|L(u)| \leq 0.2bu^2, \quad \left| L(u) + \frac{iu^3b^{3/2}}{6\sqrt{N}}A_2 \right| \leq 8\beta_4u^4/n, \quad (38)$$

where A_2 is defined as in (4). As in the proof of (6), we may obtain

$$A_2^2 = \left(N^{-1} \sum_k EV_k^3 \right)^2 \leq N^{-2} \sum_k EV_k^2 \sum_k EV_k^4 \leq N\Delta_{1n} \leq 17N\beta_4/(nb^2).$$

This together with (38) yields, for $|u| \leq \frac{1}{16}(n/\beta_4)^{1/4}$,

$$\begin{aligned} L^2(u) &\leq \left[\frac{|u|^3 b^{3/2}}{6\sqrt{N}} |A_2| + \frac{8}{n} \beta_4 u^4 \right]^2 \\ &\leq u^6 b \beta_4 / n + 128 \beta_4^2 u^8 / n^2 \leq \beta_4 (u^4 + u^6 b) / n. \end{aligned} \quad (39)$$

We are now ready to prove (25) by using (26). Rewrite (26) as

$$\exp\{u^2/2\} E \exp\{iuS_n/\sqrt{n}\} = I_1 + I_2 + I_3, \quad (40)$$

where

$$\begin{aligned} I_1 &= \sum \prod_{j=1}^n \frac{(p^j B_j)^{i_j}}{i_j!} C_{N,n,\sum_{j=1}^n j i_j}, \\ I_2 &= \sum_{i_j \geq 0, 1 \leq j \leq 3} \prod_{j=1}^3 \frac{(p^j B_j)^{i_j}}{i_j!} \left(C_{N,n,\sum_{j=1}^3 j i_j} - 1 \right), \\ I_3 &= \sum_{i_j \geq 0, 1 \leq j \leq 3} \prod_{j=1}^3 \frac{(p^j B_j)^{i_j}}{i_j!}, \end{aligned}$$

where the summation in the expression of I_1 is over all $i_j \geq 0$, $j = 1, 2, 3$ and $i_j > 0$ for at least one $j = 4, \dots, n$. As in Mirakhmedov (1983), it follows from (36)-(37) that

$$\begin{aligned} |I_1| &\leq \exp \left\{ \sum_{j=1}^3 |p^j B_j| \right\} \left(\exp \left\{ \sum_{j=4}^n |p^j B_j| \right\} - 1 \right) \\ &\leq \sum_{j=4}^n |p^j B_j| \exp \left\{ \sum_{j=1}^n |p^j B_j| \right\} \\ &\leq C n^{-1} \beta_4 u^4 \exp\{p\alpha_2 u^2/2 + 0.2bu^2\}. \end{aligned}$$

As for I_2 , it follows easily from (27) and (37) that

$$\begin{aligned} |I_2| &\leq C n^{-1} \sum_{i_j \geq 0, 1 \leq j \leq 3} \prod_{j=1}^3 \frac{|p^j B_j|^{i_j}}{i_j!} \left(\sum_{j=1}^3 j i_j \right)^2 \\ &\leq C n^{-1} \sum_{i_j \geq 0, 1 \leq j \leq 3} \prod_{j=1}^3 \frac{i_j^2 |p^j B_j|^{i_j}}{i_j!} \\ &\leq C n^{-1} \exp \left\{ \sum_{j=1}^3 |p^j B_j| \right\} \sum_{j=1}^3 (|p^j B_j| + |p^j B_j|^2) \\ &\leq C n^{-1} \beta_4 (u^2 + u^4) \exp\{p\alpha_2 u^2/2 + 0.2bu^2\}. \end{aligned}$$

We next estimate I_3 . Recalling that $b = 1 - p\alpha_2$ and noting that $I_3 = e^{\sum_{j=1}^3 p^j B_j}$, we have

$$\begin{aligned} & \left| I_3 e^{-u^2/2} - \left(1 + \frac{i^3 u^3 b^{3/2}}{6\sqrt{N}} A_2 \right) e^{-bu^2/2} \right| \\ & \leq e^{-bu^2/2} \left| e^{L(u)} - 1 - L(u) \right| + e^{-bu^2/2} \left| L(u) - \frac{i^3 u^3 b^{3/2}}{6\sqrt{N}} A_2 \right| \\ & \leq [(1/2)L^2(u)e^{|L(u)|} + 8\beta_4 u^4/n] e^{-bu^2/2} \\ & \leq C n^{-1} \beta_4 (u^4 + u^6 b) e^{-0.3bu^2}, \end{aligned}$$

where $L(u) = \sum_{j=1}^3 p^j B_j - p\alpha_2 u^2/2$ and we have used (38)-(39).

Combining (40) and all above facts for I_1 - I_3 , we obtain

$$\begin{aligned} & \left| E \exp\{iuS_n/\sqrt{n}\} - \exp\{-bu^2/2\} \left(1 + \frac{i^3 u^3 b^{3/2}}{6\sqrt{N}} A_2 \right) \right| \\ & \leq \exp\{-u^2/2\} (|I_1| + |I_2|) + \left| I_3 e^{-u^2/2} - \left(1 + \frac{i^3 u^3 b^{3/2}}{6\sqrt{N}} A_2 \right) e^{-bu^2/2} \right| \\ & \leq C n^{-1} \beta_4 (u^2 + u^4 + u^6 b) \exp\{-0.3bu^2\}, \end{aligned}$$

which implies (25). The proof of Lemma 5 is now completed. \square

After these preliminaries, we are now ready to prove the theorems.

Proof of Theorem 1. Without loss of generality, assume that $nq > 256$ and $\Delta_{1n} < 1/16$. Write $T^{-1} = \Delta_{1n} + (nq)^{-1}$ and

$$g_n(t) = \left(1 + \frac{i^3 t^3 A_2}{6\sqrt{N}} \right) \exp\{-t^2/2\},$$

where A_2 is defined as in Lemma 1. We shall prove,

(i) if $|t| \leq 1/4$, then

$$|\Psi(t) - g_n(t)| \leq C|t|T^{-1}; \quad (41)$$

(ii) if $|t| \leq \delta_0(\Delta^*)^{-1}$, where δ_0 and Δ^* are defined as in Lemma 3, then

$$|\Psi(t) - g_n(t)| \leq CT^{-1}(1+t^8)e^{-t^2/4} + C(nq)^{-3}; \quad (42)$$

(iii) if $\delta_0(\Delta^*)^{-1} \leq |t| \leq T$, then

$$|\Psi(t) - g_n(t)| \leq CT^{-1}e^{-t^2/4} + 3\sqrt{nq} \exp\{-nq\delta_N\}, \quad (43)$$

where δ_N is defined as in Theorem 1.

Note that $|A_2| \leq \sqrt{N}/4$ by $\Delta_{1n} \leq 1/16$ and the last second inequality of (39). We have $m \equiv \sup_x |G'_n(x)| \leq C(1 + N^{-1/2}|A_2|) \leq 2C$. So, by virtue of (41)-(43) and Esseen's smoothing

lemma, simple calculations show that

$$\begin{aligned} & \sup_x \left| P(S_n/\sqrt{nb} \leq x) - G_n(x) \right| \\ & \leq \left(\int_{|t| \leq 1/4} + \int_{1/4 \leq |t| \leq T_1} + \int_{T_1 \leq |t| \leq T} \frac{|\Psi(t) - g_n(t)|}{|t|} dt + CmT^{-1} \right) \\ & \leq C(\Delta_{1n} + (nq)^{-1}) + 3\sqrt{nq} \log(nq) \exp\{-nq\delta_N\}, \end{aligned}$$

where $T_1 = \min\{\delta_0(\Delta^*)^{-1}, T\}$, which implies (1) and hence Theorem 1.

We next prove (41)-(43). Throughout the proof, we write $s^2 = \psi^2 + t^2$.

Consider (42) first. Note that $g_n(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(\psi, t) d\psi$. It is readily seen that

$$|\Psi(t) - g_n(t)| \leq II_1 + II_2 + II_3 + II_4, \quad (44)$$

where

$$\begin{aligned} II_1 &= [B_n(p)]^{-1} \int_{|\psi| \leq (nq)^{1/4}/4} |h(\psi, t) - g(\psi, t)| d\psi, \\ II_2 &= [B_n(p)]^{-1} \int_{(nq)^{1/4}/4 \leq |\psi| \leq \pi\sqrt{nq}} |h(\psi, t)| d\psi, \\ II_3 &= [B_n(p)]^{-1} \int_{|\psi| \geq (nq)^{1/4}/4} \left(1 + \frac{|f(\psi, t)|}{6\sqrt{N}}\right) e^{-s^2/2} d\psi, \\ II_4 &= \left|[B_n(p)]^{-1} - (2\pi)^{-1/2}\right| \int_{-\infty}^{\infty} \left(1 + \frac{|f(\psi, t)|}{6\sqrt{N}}\right) e^{-s^2/2} d\psi. \end{aligned}$$

To estimate $II_j, j = 1, 2, 3, 4$, we first recall that, by (18), for all ψ and t ,

$$|f(\psi, t)| \leq 3s^3\sqrt{N}(\Delta_{1n} + (nq)^{-1})^{1/2} \leq \sqrt{N}s^3, \quad (45)$$

and by virtue of Stirling's formula,

$$1 \leq \sqrt{2\pi}/B_n(p) \leq 1 + 1/nq. \quad (46)$$

In view of (45) and (46), it is readily seen that

$$II_3 + II_4 \leq C(nq)^{-1}(1 + t^6)e^{-t^2/3}. \quad (47)$$

By using (22), we have

$$II_2 \leq C(nq)^{-3}. \quad (48)$$

As for II_1 , if $|t| \leq \min\{\Delta_{1n}^{-1/4}/4, \delta(\Delta^*)^{-1}\}$, Lemma 1 implies that

$$II_1 \leq C(\Delta_{1n} + (nq)^{-1})(1 + t^8)e^{-t^2/4}; \quad (49)$$

if $\Delta_{1n}^{-1/4}/4 \leq |t| \leq \delta(\Delta^*)^{-1}$, then it follows from (21) and (45) that

$$\begin{aligned} II_1 &\leq \int_{|\psi| \leq (nq)^{1/4}/4} |h(\psi, t)| d\psi + \int_{|\psi| \leq (nq)^{1/4}/4} \left(1 + \frac{|f(\psi, t)|}{6\sqrt{N}}\right) e^{-s^2/2} d\psi \\ &\leq C\Delta_{1n}e^{-t^2/4} + C_1(1 + |t|^3)e^{-t^2/2} \leq C\Delta_{1n}e^{-t^2/4}. \end{aligned} \quad (50)$$

Taking (47)-(50) into (44), we obtain the required (42).

Secondly we prove (41). Recall that $g_n(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(\psi, t) d\psi$. As in (44), we have

$$\left| \frac{d\Psi(t)}{dt} - \frac{dg_n(t)}{dt} \right| \leq III_1 + III_2 + III_3 + III_4, \quad (51)$$

where

$$\begin{aligned} III_1 &= [B_n(p)]^{-1} \int_{|\psi| \leq (nq)^{1/4}/4} \left| \frac{dh(\psi, t)}{dt} - \frac{dg(\psi, t)}{dt} \right| d\psi, \\ III_2 &= [B_n(p)]^{-1} \int_{(nq)^{1/4}/4 \leq |\psi| \leq \pi\sqrt{nq}} \left| \frac{dh(\psi, t)}{dt} \right| d\psi, \\ III_3 &= [B_n(p)]^{-1} \int_{|\psi| \geq (nq)^{1/4}/4} \left(|t| + \frac{|t||f(\psi, t)|}{6\sqrt{N}} + \frac{1}{6\sqrt{N}} \left| \frac{df(\psi, t)}{dt} \right| \right) e^{-s^2/2} d\psi, \\ III_4 &= \left| [B_n(p)]^{-1} - (2\pi)^{-1/2} \right| \int_{-\infty}^{\infty} \left(|t| + \frac{|t||f(\psi, t)|}{6\sqrt{N}} + \frac{1}{6\sqrt{N}} \left| \frac{df(\psi, t)}{dt} \right| \right) e^{-s^2/2} d\psi. \end{aligned}$$

By (18)-(19) and (46), we have that for $|t| \leq 1/4$

$$III_3 + III_4 \leq C(\Delta_{1n} + (nq)^{-1}). \quad (52)$$

By (9), (23) and (46), we have that for $|t| \leq 1/4$

$$III_1 + III_2 \leq C(\Delta_{1n} + (nq)^{-1}). \quad (53)$$

Taking these estimates into (51), we obtain for $|t| \leq 1/4$,

$$|\Psi(t) - g_n(t)| \leq |t| \sup_{|x| \leq 1/4} \left| \frac{d\Psi(x)}{dx} - \frac{dg_n(x)}{dx} \right| \leq C|t| (\Delta_{1n} + (nq)^{-1}),$$

which yields (41).

Finally we prove (43). We first notice that $\Delta_{1n} \geq 1/(16nb)$. Indeed, if $\alpha_2 \leq 1/4$, then

$$\Delta_{1n} \geq \left(N^{-1} \sum E(X_k - p\mu_k)^2 \right)^2 / (nb^2) \geq (1 - 2p\alpha_2)^2 / (nb^2) \geq 1/(16nb),$$

and if $\alpha_2 > 1/4$, then $\Delta_{1n} \geq (N^{-1} \sum \mu_k^2)^2 / (nb) = \alpha_2^2 / (nb) \geq 1/(16nb)$. This, together with the fact that

$$\Delta^* \leq \frac{9}{\sqrt{nb}^{3/2}} N^{-1} \sum_k E|X_k|^3,$$

implies that if $\delta_0(\Delta^*)^{-1} \leq |t| \leq T$, then $\delta_0 b / (9\mathcal{L}_0) \leq |t|/\sqrt{nb} \leq 16\sqrt{nb}$ and hence

$$\begin{aligned}
|h(\psi, t)|^2 &\leq \prod_k \left(1 - 2pq \left(1 - E \cos(\psi/\sqrt{nq} + tX_k/\sqrt{nb}) \right) \right) \\
&\leq \exp \left\{ -2pq \sum_k \left(1 - E \cos(\psi/\sqrt{nq} + tX_k/\sqrt{nb}) \right) \right\} \\
&\leq \exp \left\{ -2Npq \left(1 - \left| (1/N) E \sum_k \exp\{i\psi/\sqrt{nq} + itX_k/\sqrt{nb}\} \right| \right) \right\} \\
&\leq \exp \left\{ -2Npq \left(1 - \left| (1/N) \sum_k E \exp\{itX_k/\sqrt{nb}\} \right| \right) \right\} \\
&\leq \exp\{-2nq \delta_N\}. \tag{54}
\end{aligned}$$

We also note that $\Delta^* \leq 2\Delta_{1n}^{1/2}$ and this together with (45) implies that, for $\delta_0(\Delta^*)^{-1} \leq |t| \leq T$,

$$|g_n(t)| \leq \int_{-\infty}^{\infty} \left(1 + \frac{|f(\psi, t)|}{6\sqrt{N}} \right) e^{-s^2/2} d\psi \leq C(1 + |t|^3) e^{-t^2/2} \leq C\Delta_{1n} e^{-t^2/4}. \tag{55}$$

Combining (54) and (55) and using the estimate (46), we obtain that, for $\delta_0(\Delta^*)^{-1} \leq |t| \leq T$,

$$\begin{aligned}
|\Psi(t) - g_n(t)| &\leq [B_n(p)]^{-1} \int_{|\psi| \leq \pi\sqrt{nq}} |h(\psi, t)| d\psi + |g_n(t)| \\
&\leq C \Delta_{1n} e^{-t^2/4} + 3\sqrt{nq} \exp\{-nq \delta_N\},
\end{aligned}$$

which yields (43). The proof of Theorem 1 is complete. \square

Proof of Theorem 2. Without loss of generality, assume $\Delta_{2n} \leq 1$. We first prove the property (2) for $n/N \geq 1/2$ and $\Delta_{2n} \leq (nq)^{-1}/25$.

Write $T^* = (nb^2)/\beta_4$, $T_1^* = \frac{b^{1/2}}{16} \min\{(n/\beta_4)^{1/4}, \frac{1}{8}b\sqrt{n}/\mathcal{L}_0\}$ and $T_2^* = \frac{1}{15}\sqrt{nb}b^{3/2}/\mathcal{L}_0$. As in the proof of Theorem 1, it follows from Esseen's smoothing lemma that

$$\begin{aligned}
&\sup_x \left| P(S_n/\sqrt{nb} \leq x) - G_n(x) \right| \\
&\leq \left(\int_{|t| \leq T_1^*} + \int_{T_1^* \leq |t| \leq T_2^*} + \int_{T_2^* \leq |t| \leq T^*} \frac{|\Psi(t) - g_n(t)|}{|t|} dt + C \Delta_{2n} \right) \\
&= \Lambda_{1n} + \Lambda_{2n} + \Lambda_{3n} + C \Delta_{2n}, \quad \text{say.} \tag{56}
\end{aligned}$$

By virtue of Lemma 5, simple calculations show that $\Lambda_{1n} \leq C \Delta_{2n}$. Recall $\Delta_{1n} \leq 17 \Delta_{2n}$. Applying Lemma 4 and similar arguments as in the proof of (50), we obtain that $\Lambda_{2n} \leq C \Delta_{2n}$ and also $\int_{T_2^* \leq |t| \leq T^*} \frac{|g_n(t)|}{|t|} dt \leq C \Delta_{2n}$. Therefore, to prove (2), it remains to show that, for $T_2^* \leq |t| \leq T^*$,

$$|\Psi(t)| \leq 3\sqrt{n} \exp\{-n \delta_{1N}\}. \tag{57}$$

In fact, by using (3) in Zhao, Wu and Wang (2004), for $q > 0$,

$$\Psi(t) = E \exp\{itS_n/\sqrt{nb}\} = (\sqrt{2\pi}G_n(p))^{-1} \int_{-\pi}^{\pi} e^{-in\psi} \prod_k (q + pe^{i\psi + itX_j/\sqrt{nb}}) d\psi,$$

where $G_n(p) = \sqrt{2\pi}C_N^n p^n q^{N-n}$. This, together with the fact that $\delta_0 b / (5\mathcal{L}_0) \leq |t|/\sqrt{nb} \leq 16\sqrt{nb}$ whenever $T_2^* \leq |t| \leq T^*$, implies that

$$\begin{aligned} |\Psi(t)| &\leq \sqrt{2\pi}(G_n(p))^{-1} \prod_k (q + p|Ee^{itX_j/\sqrt{nb}}|) \\ &\leq \sqrt{2\pi}(G_n(p))^{-1} \exp\left\{p \sum_k (|Ee^{itX_j/\sqrt{nb}}| - 1)\right\} \\ &\leq 3\sqrt{n} \exp\{-n\delta_{1N}\}, \end{aligned}$$

where we have used the inequality $\sqrt{\pi}/2 \leq \sqrt{nq}G_n(p) < 1$ (see, for instance, Lemma 1 in Höglund(1978)). This proves (57) for $q > 0$. If $q = 0$, then $N = n$ hence by the independence of X_k ,

$$|\Psi(t)| = \prod_k |Ee^{itX_j/\sqrt{nb}}| \leq \exp\left\{\sum_k (|Ee^{itX_j/\sqrt{nb}}| - 1)\right\} \leq \exp\{-n\delta_{1N}\}.$$

This implies that (57) still holds for $q = 0$. We have now completed the proof of (57) and hence (2) for $n/N \geq 1/2$ and $\Delta_{2n} \leq (nq)^{-1}/25$.

Note that $\beta_4 \geq 1$, $\Delta_{1n} \leq 17\Delta_{2n}$ and $b \geq q \geq 1/2$ if $n/N \leq 1/2$. We have that $\Delta_{1n} + (nq)^{-1} \leq 42\Delta_{2n}$, whenever $n/N \leq 1/2$ or $\Delta_{2n} > (nq)^{-1}/25$. Based on this fact, by using a similar argument to that above and that in the proof of Theorem 1, we may obtain (2) for $n/N \leq 1/2$ or $\Delta_{2n} > (nq)^{-1}/25$, as well. The details are omitted. The proof of Theorem 2 is now complete. \square

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